

Magnetic properties and magnetic structure of the frustrated quasi-one-dimensional antiferromagnet $\text{SrCuTe}_2\text{O}_6$

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Magnetization measurements on single-crystal cubic $\text{SrCuTe}_2\text{O}_6$ with an applied magnetic field along three inequivalent high symmetry directions [100], [110], and [111] reveal weak magnetic anisotropy. The fits of the magnetic susceptibility to the result from a quantum Monte Carlo simulation on the Heisenberg spin-chain model, where the chain is formed via the dominant third-nearest-neighbor exchange interaction J_3 , yield the intrachain interaction (J_3/k_B) between 50.12(7) K for the applied field along [110] and 52.5(2) K along [100] with about the same g factor of 2.2. Single-crystal neutron diffraction unveils the transition to the magnetic ordered state as evidenced by the onset of the magnetic Bragg intensity at $T_{N1} = 5.25(9)$ K with no anomaly of the second transition at T_{N2} reported previously. Based on irreducible representation theory and magnetic space group analysis of powder and single-crystal neutron diffraction data, the magnetic structure in the Shubnikov space group $P4_132$, where the Cu^{2+} $S = 1/2$ spins antiferromagnetically align in the direction perpendicular to the spin chain, is proposed. The measured ordered moment of 0.52(6) μ_B , which represents 48% reduction from the expected value of 1 μ_B , suggests the remaining influence of frustration resulting from the J_1 and J_2 bonds.

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I. INTRODUCTION

During the past few decades, quantum magnetism in low-dimensional and frustrated systems have captured the interest of condensed-matter physicists because of their potential to exhibit exotic magnetic ground states such as spin ice [1–4], a quantum valence bond (dimer) solid [5–7], and the most sought-after quantum spin liquid [8–10]. Among these unconventional states, a quantum spin liquid has gained the most attention because its discovery and fundamental understanding could potentially yield a better understanding of other phenomena in physics such as high- T_c superconductivity [11], topological states, and anyonic physics [12], and lead to applications in quantum computing [13].

A search for this elusive quantum spin liquid has so far focused on low-dimensional and frustrated lattices, which include triangular-based lattices. Topologically, triangular [two-dimensional (2D) edge-sharing-triangle], kagome (2D

corner-sharing-triangle), hyper-kagome [three-dimensional (3D) corner-sharing triangle], and pyrochlore (3D edge-sharing-tetrahedron) lattices are considered as possible hosts of the quantum-spin-liquid ground state due to their high degree of frustration, giving rise to a macroscopically degenerate ground state, which could prompt the formation of a highly entangled quantum state. In recent years, significant progress has been made in the search, and many possible realizations of the quantum spin liquid were discovered and exhaustively studied [14]. A group of triangular-based materials that show quantum-spin-liquid traits includes $\text{ZnCu}_3(\text{OH})_6\text{Cl}_2$ [15–17], $\text{Na}_4\text{Ir}_3\text{O}_8$ [18], YbMgGaO_4 [19,20], $\text{Ca}_{10}\text{Cr}_7\text{O}_{28}$ [21,22], $\text{Ce}_2\text{Zr}_2\text{O}_7$ [23], and $\text{PbCuTe}_2\text{O}_6$ (Cu^{2+} , $S = 1/2$) [24,25].

Studies of $\text{PbCuTe}_2\text{O}_6$ revealed the absence of a long-range λ -like transition down to 2 K and possible emergence of quantum spin liquid at low temperatures [24–26]. On the other hand, $\text{SrCuTe}_2\text{O}_6$ with an almost identical crystal structure exhibits two successive magnetic phase transitions, one of which is to a magnetically ordered state at the Néel temperature T_{N1} of 5.5 K [27,28]. In order to decipher the underlying mechanism that gives rise to different ground states in these two

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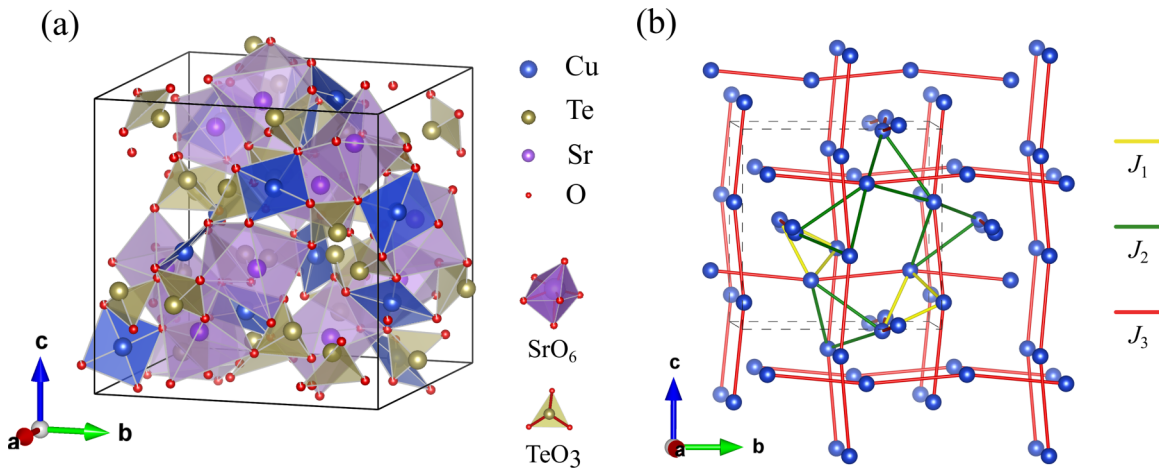


FIG. 1. (a) The crystal structure of $\text{SrCuTe}_2\text{O}_6$ consists of CuO_4 square plaquettes (blue), TeO_3 triangular plaquettes (yellow), and SrO_6 octahedra (purple). (b) The intertwined spin network of Cu^{2+} $S = 1/2$ spins is connected by the exchange interactions J_1 (yellow), J_2 (green), and J_3 (red) to form isolated triangles, a hyper-kagome lattice, and spin chains, respectively.

seemingly similar systems, a detailed study on a single-crystal sample is required. Hence, in this article we report the studies of magnetic properties and a magnetic structure on powder and single-crystal $\text{SrCuTe}_2\text{O}_6$.

Noncentrosymmetric $\text{SrCuTe}_2\text{O}_6$ crystallizes in the chiral cubic space group $P4_132$ (No. 213) with lattice parameter $a = 12.473(1)$ Å [29]. The crystal structure consists of CuO_4 square plates, units of TeO_3 , and oxygen octahedra surrounding strontium (Sr) atoms as shown in Fig. 1(a). The Cu^{2+} ions in the CuO_4 plaquettes carry spin $1/2$ and give rise to the magnetic properties of this system. As with $\text{PbCuTe}_2\text{O}_6$, the spin network of $\text{SrCuTe}_2\text{O}_6$ consists of three types of intertwined networks connected by the nearest-, second-nearest-, and third-nearest-neighbor exchange interactions J_1 , J_2 , and J_3 , respectively [27,28], as shown in Fig. 1(b). The Cu^{2+} ions connected by J_1 form an isolated triangle with Cu-O-Sr-O-Cu superexchange pathways. The second-nearest-neighbor network formed by J_2 via Cu-O-Te-O-Cu pathways connect the Cu^{2+} ions to form a hyper-kagome network. Finally, J_3 links the Cu^{2+} ions to form spin chains along the cubic crystallographic axes. First principle calculations performed to estimate the strength of these exchange interactions showed that for $\text{SrCuTe}_2\text{O}_6$, J_3 is the most dominant with J_2 about 10% of J_3 and J_1 about 1% of J_3 , and established that $\text{SrCuTe}_2\text{O}_6$ is a spin-chain system with relatively weak and intricate interchain couplings [27]. In contrast, for $\text{PbCuTe}_2\text{O}_6$, J_2 is the most dominant interaction, and hence this system can be characterized as the 3D hyper-kagome antiferromagnet, in which geometric frustration can potentially suppress the Néel state and prompt the emergence of the quantum spin state [24–26].

Previous magnetic susceptibility and heat capacity measurements on a $\text{SrCuTe}_2\text{O}_6$ powder sample [27,28,30] revealed magnetic transitions at $T_{\text{N1}} = 5.5$ K, which is characterized as the paramagnetic-to-antiferromagnetic transition, and an unexplained anomaly at T_{N2} , which is smaller than T_{N1} and field dependent [27,28]. The magnetic susceptibility data also showed a broad maximum, which is typical in a spin-chain system, at around 32 K suggesting short-range correlation. A fit of the high-temperature data to the Curie-

Weiss law yields the Curie-Weiss temperature $\Theta_{\text{CW}} = -35$ K. Assuming that the spin-chain network is connected by J_3 , the dominant exchange interaction (J_3/k_B) is estimated to be 49 K as compared to the first-principle-calculation value of 45 K [27]. Furthermore, multiple magnetic transitions with a nontrivial $H - T$ phase diagram were observed for magnetizations up to 9 T [27,28].

In this work, magnetic properties of $\text{SrCuTe}_2\text{O}_6$ are studied using magnetization measurements on a single-crystal sample, and the magnetic transition at T_{N1} and magnetic structure of the ordered state are investigated using neutron diffraction. This article is organized as follows. In Sec. II, the single-crystal synthesis and experimental methods are described. The result of the single-crystal x-ray diffraction is discussed in Sec. III A, macroscopic magnetic properties on the single crystals are investigated and analyzed using quantum Monte Carlo simulations in Sec. III B, and powder and single-crystal neutron diffraction data are discussed in Sec. III C. The article ends with the conclusion in Sec. IV.

II. EXPERIMENTAL DETAILS

A powder sample of $\text{SrCuTe}_2\text{O}_6$ was first synthesized by the standard solid-state reaction of high purity SrCO_3 , TeO_2 , and CuO . The preparation method is described elsewhere [28]. The obtained pure phase powder of $\text{SrCuTe}_2\text{O}_6$ was then used as a starting material for single-crystal growth using the vertical gradient freeze technique. The powder was loaded into a pointed-bottom alumina tube (recrystallized alumina 99.8%). The crucible was then sealed in an evacuated quartz tube, which is crucial to minimize the formation of $\text{SrCuTe}_2\text{O}_7$ as an impurity phase. The sample was melted at 800°C and held at this temperature for 24 h to ensure a homogeneous and complete melt, before subjecting the sample to a $20^\circ\text{C}/\text{cm}$ temperature gradient in a vertical furnace at a rate of $1\text{ cm}/\text{day}$. After the sample reached the position with $T = 650^\circ\text{C}$, the furnace was cooled to room temperature at a rate of $200^\circ\text{C}/\text{h}$. The sample was then mechanically extracted from the crucible. Single crystals with the largest mass of about 1 g were obtained.

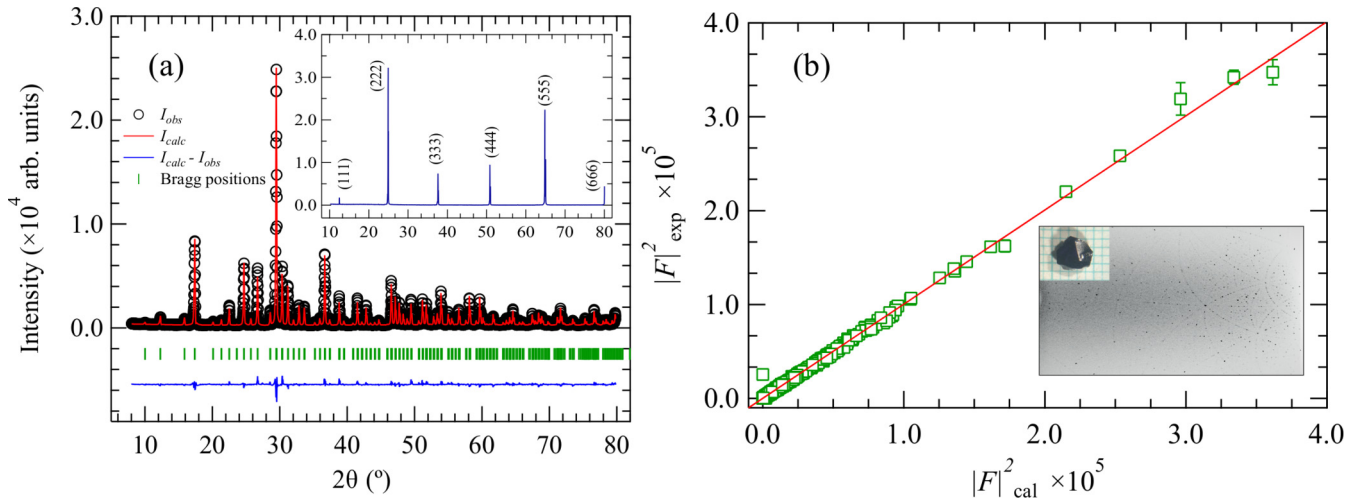


FIG. 2. (a) Powder x-ray diffraction pattern measured on the crushed single crystals of $\text{SrCuTe}_2\text{O}_6$ at room temperature. The black circle denotes the observed data, the red line the calculated pattern, the vertical green line the Bragg positions, and the blue line the difference between the observed and calculated patterns. The inset shows the θ - 2θ scan on the cleaved (1,1,1) facet of single-crystal $\text{SrCuTe}_2\text{O}_6$. (b) The refinement result of the single-crystal x-ray diffraction data shows the agreement between the measured and calculated scattering intensities. Error bars represent one standard deviation. The inset shows a photograph of a crystal and Laue x-ray diffraction measured on the cleaved surface.

Small pieces of the crystals were collected and ground thoroughly for a powder x-ray diffraction measurement to confirm sample purity. The powder x-ray diffraction data was fitted using the Rietveld method implemented in FULLPROF [31]. Single-crystal x-ray diffraction was also performed using a Bruker X8 APEX II CCD Diffractometer with $\text{MoK}\alpha$ radiation at room temperature. The refinement on the single-crystal diffraction data for fractional coordinates was done using SHELXL [32].

In order to investigate the macroscopic magnetic properties of $\text{SrCuTe}_2\text{O}_6$ in the single-crystal sample, the magnetic susceptibility was measured with the applied field aligned along the three inequivalent directions [100], [110], and [111]. A single crystal was aligned using a four-circle x-ray diffractometer and cut into a cube with dimensions of $1 \times 1 \times 1 \text{ mm}^3$ ($\sim 20 \text{ mg}$). The aligned crystal was then attached to a Teflon rod using GE-7031 varnish and placed inside a measuring stick. The magnetic susceptibility was measured as a function of temperature from 2 K up to 300 K with the applied magnetic field of 1.0 T using a superconducting quantum interference device (MPMS-XL, Quantum Design). To analyze the susceptibility data, quantum Monte Carlo (QMC) simulations were performed to capture the broad maximum of the magnetic susceptibility data in order to extract the value of the leading exchange interaction. The QMC simulations were performed using the LOOP algorithm [33] in the ALPS simulation package [34] on a cluster of 100 spins for the spin-chain model and 96 000 spins ($20 \times 20 \times 20$ unit cells) for the J_3 model of $\text{SrCuTe}_2\text{O}_6$ in the temperature range of $0.01 \leq t \leq 10$ ($t = k_B T/J$) with 100 000 Monte Carlo steps for thermalization and 500 000 Monte Carlo steps after thermalization. The numerical result was fitted with the experimental data by a Padé approximant [35] and the leading exchange interaction and Landé g factor were finally obtained.

To study the magnetic structure, powder neutron diffraction was performed at the high-resolution neutron powder

diffractometer BT1 at the NIST Center for Neutron Research (NCNR), USA. The $\text{Ge}(311)$ monochromator was used to select neutrons with $\lambda = 2.077 \text{ \AA}$ and collimations of $60' - 20' - 7'$ were employed. In addition, elastic neutron scattering on a small 130-mg single crystal was conducted using the Double Focusing Triple-Axis Spectrometer BT7 [36] at NCNR. Neutron scattering measurement was performed using fixed incident energy of 14.7 meV and pyrolytic graphite $\text{PG}(002)$ was used as a monochromator. The sample was cooled to the base temperature using a closed cycle ^4He cryostat. A position sensitive detector (PSD) was employed in the two-axis mode with collimations of open $- 80' - 80' - \text{PSD}$. The sample was aligned so that the hhl plane was in the scattering plane.

III. RESULTS AND DISCUSSION

A. X-ray diffraction

The x-ray diffraction data, which were measured on the powder sample obtained by grinding small pieces of the single crystals, along with the result of the Rietveld refinement are shown in Fig. 2(a). The results show that the obtained $\text{SrCuTe}_2\text{O}_6$ single crystals are single phase without any trace of impurity. However, we will later show that from the neutron diffraction data measured on an as-grown powder sample, $\text{SrCuTe}_2\text{O}_7$ is present as an impurity. The absence of this impurity phase in the single-crystal x-ray diffraction is due to the fact that only small pieces of the $\text{SrCuTe}_2\text{O}_6$ single crystals were selected for grinding, and hence $\text{SrCuTe}_2\text{O}_7$, which also showed up as having a different color, was selected out. The lattice parameter obtained from the refinement is $a = 12.466(6) \text{ \AA}$ for the cubic space group $P4_132$ in agreement with previous reports [27–30]. The residual parameters for the Rietveld refinement are $R_p = 6.66\%$, $R_{wp} = 8.49\%$, and the goodness of fit (GoF) of 2.1. A θ - 2θ scan shown in the inset of Fig. 2(a) on the naturally cleaved facet confirms single

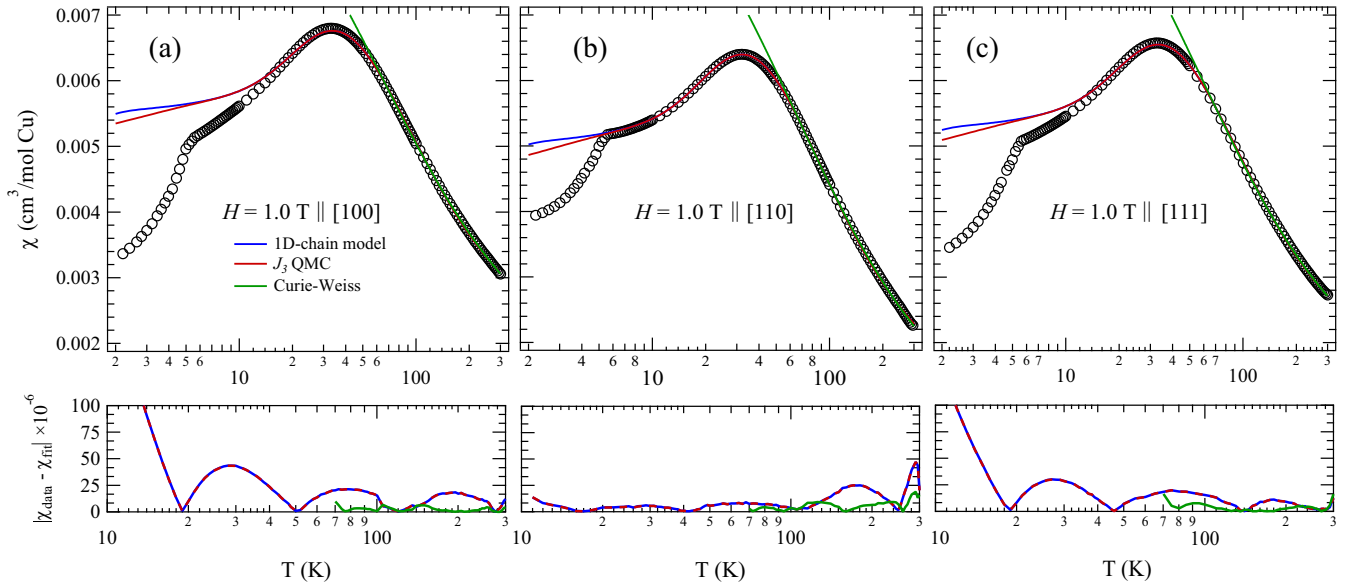


FIG. 3. Magnetic susceptibility measured on a single crystal of $\text{SrCuTe}_2\text{O}_6$ with applied magnetic fields of 1.0 T along three crystallographic directions: (a) [100], (b) [110], and (c) [111]. The blue and red lines represent the fits to the 1D-chain model and our QMC simulations, respectively. The green line denotes the Curie-Weiss law fit. The lower panels show the residuals between the data and the fits.

crystallinity and shows that the cleaved facet is the 111 plane. To further investigate the quality of the single crystals, Laue photography was also performed in the transmission mode. The result shown in the inset of Fig. 2(b) reveals clear Bragg peaks confirming the quality of the crystal.

Single-crystal x-ray diffraction at room temperature was also performed and the data were refined to extract the atomic fractional coordinates. The refinement was done on 1192 unique reflections with $F_{\text{obs}} > 4\sigma(F_{\text{obs}})$ yielding the fitted lattice parameter $a = 12.472(6) \text{ \AA}$. Figure 2(b) shows the agreement between the measured and calculated intensity resulting from the refinement with the residual parameters $R_1 = 2.47\%$, $wR_2 = 6.01\%$, and $\text{GoF} = 1.042$. We note that R_1 for the refinement performed on all (1263) reflections is 2.79%. The fractional coordinates of $\text{SrCuTe}_2\text{O}_6$ obtained from the refinements of the powder (the crushed crystalline sample) and single-crystal x-ray diffraction data are summarized in Appendix A.

B. Magnetic susceptibility

Figure 3 shows the magnetic susceptibility measured with the applied field along three inequivalent directions [100], [110], and [111]. The susceptibility data are plotted against the temperature in a log scale to emphasize the broad maximum at $\sim 30 \text{ K}$, and the kink at the Néel temperature $T_{\text{N1}} = 5.25(5) \text{ K}$ (T_{N1} was obtained from the neutron scattering data shown in Fig. 4). Below T_{N1} , the magnetic susceptibility decreases, which is indicative of the antiferromagnetic arrangement of the spins. The broad maximum is a sign of short-range spin correlation, typical for low-dimensional antiferromagnetic systems [37], and can be used to estimate the leading exchange interactions of the system. This broad peak was consistent with that previously observed in the powder sample [27,28]. We note that for the data measured in the applied field of 1.0 T on the powder sample in Refs. [27,28], the anomaly at

T_{N2} is absent, possibly because it is very weak at low field; for the magnetization measurements, this anomaly, which is field dependent, appears below T_{N1} for $H \geq 2.0 \text{ T}$ [27]. Consistent with the previous work on the powder sample, we were able to observe the second magnetic transition at higher magnetic fields (not shown here).

To estimate the dominant exchange coupling between the Cu^{2+} $S = 1/2$ spins and verify the connectivity of the spin network, we fit the high-temperature ($T \gtrsim T_{\text{N1}}$) data to the

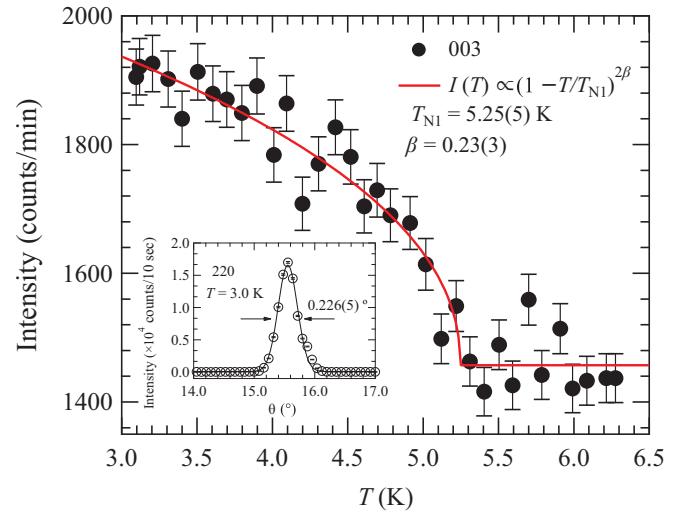


FIG. 4. Neutron diffraction intensity of the 003 reflection as a function of temperature. The red line denotes a power-law fit for the magnetic scattering intensity representing the order parameter. Error bars represent one standard deviation. The inset shows the θ scan of the 220 structural reflection. We note that error bars are smaller than the plot symbol. The solid line represents the Gaussian fit to capture the line shape of the peak.

TABLE I. The fitted parameters of magnetic susceptibilities data along [100], [110], and [111] with the mean-field Curie-Weiss law [Eq. (1)].

$H \parallel$	χ_0 (cm ³ /mol Cu) (fixed)	Θ_{CW} (K)	μ_{eff} (μ_B)
[100]	1.69×10^{-3}	$-46.5(6)$	$1.97(1)$
[110]	0.834×10^{-3}	$-41.6(8)$	$2.01(2)$
[111]	1.31×10^{-3}	$-44.3(6)$	$1.99(1)$

mean-field and QMC calculations. The magnetic susceptibility data above 70 K along the three axes were first fitted with the mean-field Curie-Weiss law,

$$\chi(T) = \chi_0 + \frac{C}{T - \Theta_{\text{CW}}}, \quad (1)$$

where C and Θ_{CW} are the Curie constant and Curie-Weiss temperature, respectively. χ_0 , which is fixed to the values obtained from the fits of the QMC results (later discussed), represents the temperature-independent background, which is relatively high (about half of the measured susceptibility at high temperatures) due to the Teflon used to fix the crystal. We note that χ_0 is positive and its magnitude is about an order of magnitude larger than the core diamagnetic and Van Vleck paramagnetic susceptibilities [27]. Teflon is diamagnetic and hence gives rise to negative magnetic susceptibility. However, in the measurements, the signal from the Teflon rods was treated as background by MPMS. As a result, the empty gap between the Teflon rods gives rise to positive magnetic-susceptibility background. The difference in the value of χ_0 could be due to the difference in the gap between the two Teflon rods that were used to fix the aligned crystal in different orientations. The Curie-Weiss fitted curves are shown by the green lines in Figs. 3(a)–3(c). The fitted parameters along each field direction are shown in Table I. The effective magnetic moment, $\mu_{\text{eff}} = \sqrt{3k_B C/N_A}$, obtained for the three inequivalent directions does not significantly deviate from one another, which is indicative of relatively weak spin anisotropy with a similar g factor. All are found to be slightly higher but still close to the spin-only value of $\mu_{\text{eff}} = g\mu_B \sqrt{S(S+1)} = 1.73 \mu_B$ for $g = 2$ and $S = 1/2$. The slight increase of the value of μ_{eff} from the spin-only value could indicate the presence of the spin-orbit coupling and the orbital contribution to the magnetic moment. For $S = 1/2$, the measured μ_{eff} of 2.0 implies $g = 2.3$. The obtained Curie-Weiss temperature, which is close to the previous reported values measured on the powder sample [27,28], as shown in Table I, indicates no significant deviation along the three field directions. The negative Curie-Weiss temperature indicates that the dominant exchange interaction is antiferromagnetic. The order of frustration $f = |\Theta_{\text{CW}}/T_{\text{N1}}|$ for $T_{\text{N1}} = 5.25$ K and the average Curie-Weiss temperature of -44 K, yields $f \sim 8$ suggesting some degree of frustration in this system that could result from J_1 and J_2 that give rise to the frustrated spin networks, isolated triangles, and the hyper-kagome lattice, respectively.

The decrease of the magnetic susceptibility for $T < T_{\text{N1}}$ suggests an antiferromagnetically ordered state with the absence of ferromagnetism, which in some materials could be due to spin canting. The absence of the spin canting is

consistent with the absence of the staggered Dzyaloshinskii-Moriya (DM) interactions, which can give rise to canted moments, between two spins connected by J_3 . However, the chiral $P4_132$ symmetry of the underlying crystal structure can lead to the uniform DM interactions, which, in combination with the complex spin network, can result in the nontrivial $H - T$ phase diagram observed in previous work [27]. It was discovered that a chiral material with space group $P4_132$ can host magnetic skyrmions [38], the existence of which can be explained by the DM interaction. In addition, the similar drop of the magnetic susceptibility for $T \lesssim T_{\text{N1}}$ along all three field directions [Figs. 3(a)–3(c)] further suggests that there is no global easy axis and the antiferromagnetic alignment must be along a local easy axis determined by the local environment around the Cu^{2+} ions.

To go beyond the mean-field approximation, we performed a QMC simulation to calculate the magnetic susceptibility that was subsequently used to fit the experimental data. The previous work showed that the broad maximum observed in $\text{SrCuTe}_2\text{O}_6$ can be well captured by the one-dimensional (1D) spin-chain model [27,28]. For this work, we performed QMC simulations with the LOOP algorithm based on the one-dimensional spin-chain model and on the spin network connected by J_3 [shown in Fig. 1(b)]. We note that the two models without any interchain coupling are in fact equivalent. However, the J_3 spin network was used for possible inclusion of J_1 and J_2 , which serve as interchain interactions. The simulated result was then fitted to the experimental data using

$$\chi(T) = \chi_0 + \chi_{\text{QMC}}(T), \quad (2)$$

where

$$\chi_{\text{QMC}}(T) = \frac{N_A \mu_B^2 g^2}{k_B J} \chi^*(k_B T/J). \quad (3)$$

N_A , μ_B , and k_B are the Avogadro constant, Bohr magneton, and Boltzmann constant, respectively. J and g are fitted to the experimental data. The susceptibility $\chi^*(t)$ as a function of the reduced temperature $t = k_B T/J$ was obtained by fitting the QMC results using a Padé approximant [35,39]. The fitting was performed on the experimental data from $T = 300$ K down to 10 K, which is slightly above T_{N1} and contains the broad peak, along the three field directions as shown in Fig. 3. The QMC result appears to fit the data for the applied field along the [110] direction very well, where the QMC calculations and the experimental data are in good agreement around the broad peak and the consistency extends down to the temperature just above T_{N1} . In contrast, for the [100] and [111] data, the QMC results appear to deviate from the data around the broad peak. We cannot explain why the result of the QMC calculations fits the [110] data much better than it does the other two field directions. The fitted parameters are summarized in Table II. The fitted values of the Landé g factor are consistent with those obtained from the Curie-Weiss fit discussed in the previous section. It can be seen that the obtained exchange interactions along three axes do not deviate much and the values are very close to those obtained from the mean-field approximation. This suggests that the 1D spin-chain model along third-nearest neighbors adequately describes the macroscopic magnetic properties. The obtained Landé g factor along the three inequivalent field directions

TABLE II. The fitted parameters of the magnetic susceptibilities data along [100], [110], and [111] with the QMC simulation.

$H \parallel$	χ_0 (cm ³ /mol Cu)	J/k_B (K)	g factor
[100]	$1.69(1) \times 10^{-3}$	52.5(2)	2.197(6)
[110]	$0.834(5) \times 10^{-3}$	50.12(7)	2.249(2)
[111]	$1.31(1) \times 10^{-3}$	51.8(1)	2.221(5)

from the QMC fit are also consistent with the spin network where the chains are running along the three crystallographic axes giving rise to the relatively isotropic magnetic susceptibility in this cubic system. We attempted to include J_1 and J_2 in our QMC calculations but encountered the sign problem due to the frustration of the J_1 and J_2 bonds. As a result, we were unable to obtain reliable J_1 and J_2 values especially around the broad peak and the low-temperature region.

C. Single-crystal and powder neutron diffraction

To investigate the microscopically magnetic properties of SrCuTe₂O₆, elastic neutron scattering was performed on a small crystal (≈ 130 mg) at BT7. The inset of Fig. 4 shows a θ scan (rotating only the sample) around the 220 structural Bragg peak with a Gaussian fit yielding a full-width-at-half-maximum (FWHM) equal to $0.226(5)^\circ$, indicative of good crystallinity. The scattering intensity of the 003 and 113 reflections was measured as a function of temperature. The results show the onset of the intensity increase at T_{N1} . The extra scattering intensity below T_{N1} is indicative of magnetic scattering resulting from the ordering of the magnetic moments. Figure 4 shows the 003 scattering intensity data, which also serve as a measure of the order parameter below T_{N1} , along with the power law fit. The fit to $I(T) \propto (1 - T/T_{N1})^{2\beta}$ for the order parameter yields the critical exponent $\beta = 0.23(3)$ and Néel temperature $T_{N1} = 5.25(5)$ K, where the errors represent one standard deviation. We note that the magnetic scattering intensity is proportional to M^2 , where M is the sublattice magnetic moment, and hence the factor of two in the exponent. For the 113 data (not shown), the critical exponent and Néel temperature are $0.27(4)$ and $5.12(5)$ K, respectively, which are consistent with those obtained from the 003 data. The fitted value of β is typical for low-dimensional magnetic systems [40–42], the 1D spin chain in this case. The obtained value of T_{N1} is consistent with magnetic susceptibility data and with previous reports on the powder sample [27,28,30]. However, given the resolution of the data, the order parameter in Fig. 4, which was measured at zero magnetic field, does not show an anomaly of the second magnetic transition around T_{N2} . It is possible that the second transition can only be detectable at high field where the anomaly becomes stronger as suggested by the magnetization data. Hence in order to investigate the magnetic structure change at T_{N2} , future in-field elastic neutron scattering is required.

To determine the magnetic structure of SrCuTe₂O₆, powder neutron diffraction was performed at BT1 and the data were collected at 10 K and 1.5 K, above and below T_{N1} , respectively. The refinement of the nuclear structure was first performed on the 10-K data with the fitted lattice parameter $a = 12.4321(5)$ Å. The resulting Rietveld refinement

[the inset of Fig. 5(a)], which yields $R_p = 3.72\%$, shows results that are consistent with the crystal structure of SrCuTe₂O₆ obtained from the x-ray diffraction data (Appendix A). The sample, however, contained some impurities, the majority of which was identified to be SrCuTe₂O₇ that constitutes roughly 2.1 wt.%. The proximity of impurity reflections to some of the magnetic Bragg reflections and weak magnetic intensity hinder the refinement of the magnetic structure from the powder neutron diffraction data. As a result, the fitted ordered magnetic moment has large error as will be discussed below.

The magnetic structure of SrCuTe₂O₆ was first analyzed using the irreducible representation theory. The detail was described in Appendix B. We note that due to the large number of free parameters, we were unable to perform full refinement for Γ_3 , Γ_4 , and Γ_5 with 6, 12, and 15 free parameters, respectively. Therefore, we have to rely on the magnetic space group analysis in order to further subclassify possible magnetic structures of SrCuTe₂O₆ and reduce the number of free parameters.

Based on the Landau-type transition with a single order parameter, the magnetic Shubnikov space groups can be derived from the paramagnetic parent space group $P4_1321'$ giving rise to 14 Shubnikov space groups as shown by the graph of subgroups, which was generated using *k*-SUBGROUPSMAG [43], in Fig. 6. Out of these 14 subgroups, there are a total of five maximal magnetic subgroups, $P4_132$ (No. 213.63), $P4_1'32'$ (No. 213.65), $P4_12_1'2'$ (No. 92.114), $P4_1'2_1'2$ (No. 92.115), and $C22'2_1'$ (No. 20.34). $P4_132$ and $P4_1'32'$ correspond to Γ_1 and Γ_2 , respectively, whereas $P4_12_1'2'$, $P4_1'2_1'2$, and $C22'2_1'$ correspond to Γ_4 and Γ_5 . Since some of the basis vectors of Γ_4 and Γ_5 are absent for the magnetic space groups $P4_12_1'2'$, $P4_1'2_1'2$, and $C22'2_1'$, the number of fitting parameters is reduced.

Assuming that symmetry reduction at the magnetic transition to the ordered state is minimal, we performed the refinement of the magnetic structure on the 1.5-K data for the four (out of five) maximal subgroups, namely $P4_132$ (Γ_1) with one free parameter, $P4_1'32'$ (Γ_2) with two free parameters, $P4_12_1'2'$ with five free parameters, and $P4_1'2_1'2$ with four free parameters. We note that an attempt to perform the refinement for the maximal subgroup $C22'2_1'$ with nine free parameters was unsuccessful. All fitting parameters including the lattice parameter, atomic positions, and peak profile parameters, were kept constant and the same as those obtained from the fitting of the 10-K data. However, the background was adjusted due to the difference in the incident neutron flux. Since the magnetic scattering was observed at low momentum transfer, and hence low 2θ angles, the magnetic-structure refinement was performed for $20.5^\circ < 2\theta < 22.5^\circ$ and $28.5^\circ < 2\theta < 32.5^\circ$, where four magnetic Bragg reflections were observed. Figure 5(a) shows the difference pattern between 1.5 K and 10 K for these four magnetic reflections, 012, 003, 013, and 113, two of which [003 and 013] appear next to the impurity peaks (not shown). The figure also shows that the calculated intensity based on the Shubnikov space group $P4_132$, which yields the R factor of 5.51%, fits the data slightly better than that based on $P4_1'32'$ with the R factor of 7.02%. The fitted magnetic moment for $P4_132$ is $0.8(7) \mu_B$, where the error was estimated from fitting the 10-K data; the large error is

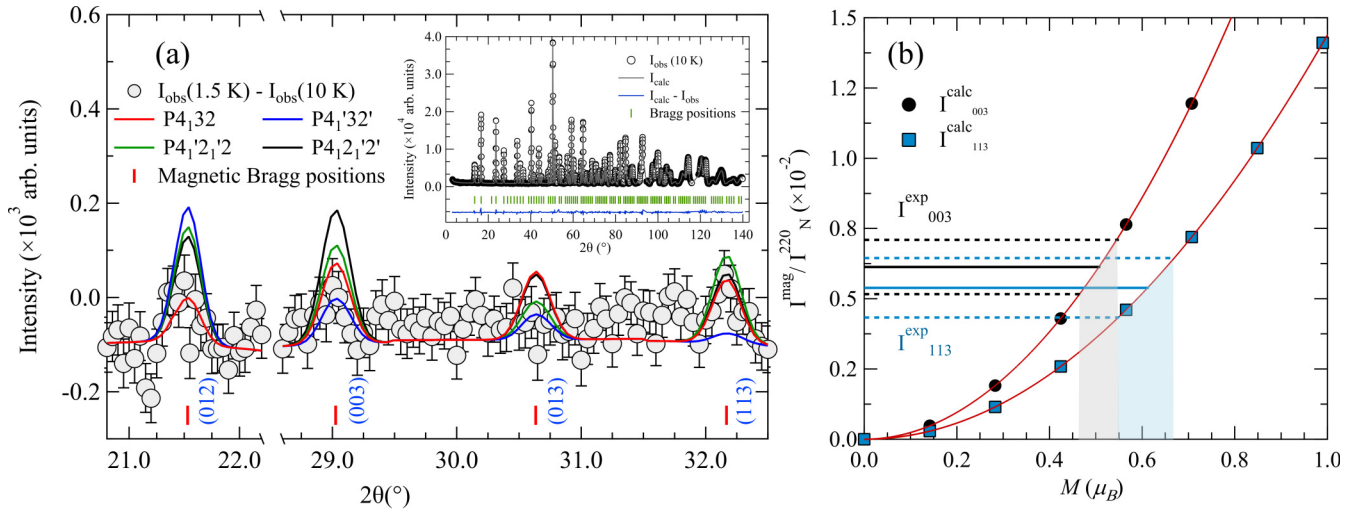


FIG. 5. (a) The difference data between 1.5 K and 10 K show the magnetic Bragg scattering at 012, 003, 013, and 113. The lines represent the magnetic-structure refinements based on magnetic space groups $P4_1'32'$, $P4_132$, $P4_12_1'2'$, and $P4_1'2_1'2$. The vertical red lines denote the magnetic Bragg reflections. The powder neutron diffraction data were measured at 10 K (the inset). The black symbols denote the observed data, the grey line the calculated nuclear-scattering pattern, the vertical green lines the Bragg positions, and the blue line the difference between the observed and calculated data. (b) The calculated single-crystal magnetic intensities for the magnetic space group $P4_132$ are shown as a function of ordered moment (M) for the reflections 003 (black circles) and 113 (blue squares). The red lines denote the fits to the quadratic function. The black (blue) horizontal line represents the experimental value of the ratio between the magnetic scattering intensity measured at 003 and 113, and the structural scattering intensity measured at 220. The gray (blue) shaded region indicates a range of the measured ordered moment for 003 and 113.

due to the significant contribution from the nearby impurity peaks and weakness of the magnetic signal. The refinement based on $P4_12_1'2'$ and $P4_1'2_1'2$ yields the R factors of 6.31% and 6.38%, respectively, which are marginally worse than that for $P4_132$. We note that even though the number of fitting parameters for $P4_1'32'$, $P4_12_1'2'$, and $P4_1'2_1'2$ is higher than that

for $P4_132$, the fitted results for $P4_1'32'$, $P4_12_1'2'$, and $P4_1'2_1'2$ do not become better. However, the small difference of the R factors is not statistically significant enough to validate that $\text{SrCuTe}_2\text{O}_6$ magnetically orders in the magnetic space group $P4_132$. Furthermore, we were unable to rule out $C22_1'2_1'$, nor, under the assumption of minimal symmetry reduction at

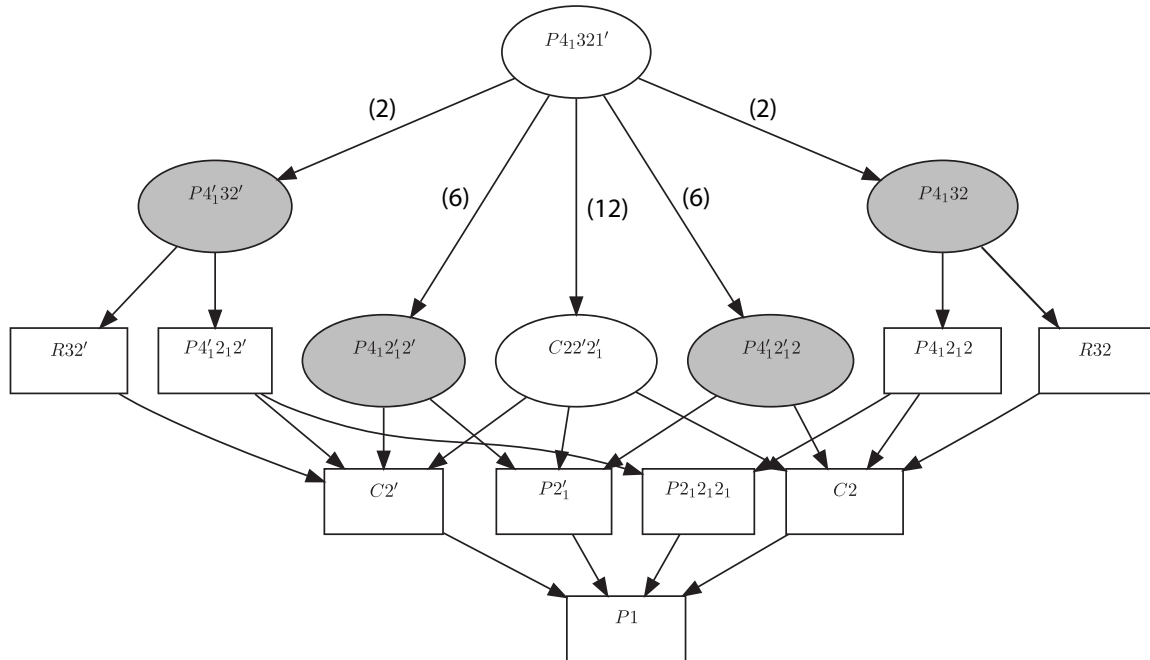


FIG. 6. A diagram of subgroups shows a hierarchy of possible subgroups of the paramagnetic parent space group $P4_1321'$. The maximal subgroups are indicated by eclipses, and those that were used in the refinement are highlighted by shading. A subgroup index between the parent space group and a maximal subgroup is shown in parentheses. The diagram is generated using k -SUBGROUPSMAG [43].

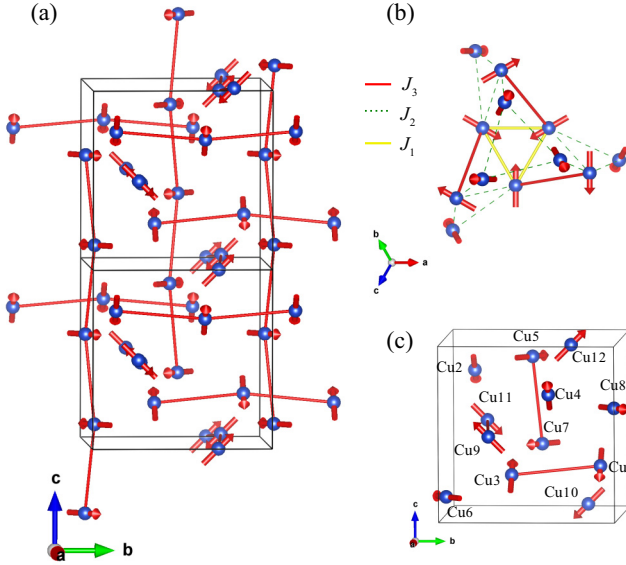


FIG. 7. (a) The magnetic structure of SrCuTe₂O₆ belongs to the Shubnikov space group $P4_132$ (Γ_1). The dominant J_3 antiferromagnetic exchange interactions form spin chains along the crystallographic axes. (b) The spin network, which is formed by J_1 (yellow), J_2 (dashed green), and J_3 (blue), is shown along with the spin structure. (c) The Cu atoms in a unit cell are labeled according to Table IV.

the transition, examine nonmaximal subgroups. As previously noted, since the magnetic scattering in SrCuTe₂O₆ is weak and some magnetic Bragg reflections are in close proximity to impurity peaks, the refinement of the powder neutron diffraction data yields an inconclusive result with large error. Hence, in this work, $P4_132$ is proposed as the most likely candidate based on the magnetic space group analysis.

The resulting magnetic structure for $P4_132$ is shown in Fig. 7. The magnetic moments of the Cu²⁺ $S = 1/2$ spins in SrCuTe₂O₆ antiferromagnetically align in the direction perpendicular to the chain formed by J_3 , which is the most dominant exchange interaction, consistent with proposed spin network deduced from the DFT calculations [27]. This antiferromagnetic spin structure is consistent with the magnetization data (discussed above) where no weak ferromagnetism, which could result from spin canting, was observed. Interestingly, we observed that the spins on the corners of an isolated triangle connected by J_1 , i.e., the weakest exchange interaction among the three considered, form a co-planar 120° configuration [Fig. 7(b)], which relieves, to some degree, the geometrical frustration inherent in the triangle-based spin network. However, the J_2 interactions, which form the hyper-kagome spin network, remain highly frustrated. The magnetic order in this quasi-1D system is most likely stabilized by this intricate network of further-nearest-neighbor interactions [Figs. 1(b) and 7(b)].

In order to better extract the value of the ordered moment, we performed the analysis on the single-crystal neutron diffraction data measured at BT7. In Fig. 4, we were able to clearly observe the magnetic Bragg intensities at 003 and 113 (not shown), which we will denote as I_M^{003} and I_M^{113} , respectively. The base temperature of 3 K for the single-

crystal experiment might not be low enough relative to T_N to give a good estimate of the ordered moment as suggested by the increasing trend of the scattering intensity in Fig. 4. Hence, I_M^{003} and I_M^{113} were obtained by extrapolating the order parameter fitted curve of the scattering intensity measured at 003 and 113, respectively, to 1.5 K, at which the power neutron diffraction was measured. In comparison with the nuclear Bragg intensity at 220, I_N^{220} , we calculated the ratio between the magnetic and nuclear scattering intensity as $I_M^{003}/I_N^{220} = 6.1(10) \times 10^{-3}$ and $I_M^{113}/I_N^{220} = 5.4(11) \times 10^{-3}$. These obtained values of the magnetic to nuclear intensity ratio are very weak, i.e., roughly of the same order of magnitude as the statistical error in the powder data, and hence small contributions from the nearby impurity scattering can cause a large error in the refinement. As a result, the magnetic scattering is barely noticeable in the difference plot between the 1.5-K and 10-K data [Fig. 5(a)], and the magnetic structure refinement on the powder neutron diffraction data fails to yield a reliable result.

To obtain the value of the ordered magnetic moment from the single-crystal data, we compare the magnetic intensities I_M^{003}/I_N^{220} and I_M^{113}/I_N^{220} from the single-crystal data and the magnetic scattering intensity calculated from FULLPROF for $P4_132$. We convert the integrated intensities calculated from FULLPROF for powder to those for single-crystal by multiplying $\sin \theta / m_{hkl}$, where m_{hkl} is the multiplicity of the hkl reflection [44]. The result is shown in Fig. 5(b); the red curves in Fig. 5 denote a fit to $I^{\text{mag}}/I^{220} \propto M^2$. We note that M is the only free parameter for $P4_132$. Given the values of I_M^{003}/I_N^{220} and I_M^{113}/I_N^{220} from above, we estimate the ordered moment to be $0.51(4)\mu_B$ and $0.61(6)\mu_B$, respectively. The horizontal black and blue solid lines denote the values of the intensity ratios for 003 and 113, respectively, with the dashed lines representing the range of the error. We have done a similar analysis for $P4_132'$, which has two free parameters, C_1 and C_2 . We found that if C_2 is equal to zero, the calculated magnetic intensity of 113 will be greater than that of 003, which is inconsistent with the experimental data. With increasing C_2 , the 003 intensity can become larger than the 113 intensity but they are still inconsistent with the experimental data. In addition, refinement of the single-crystal data was also performed using JANA2006 [45] for magnetic space groups $P4_132$. We were unable to check $P4_132'$, $P4_12_1'2_1'$, and $P4_12_1'2_1'$ due to a limited number of data points. The least-square refinement yields $wR_2 = 4.63\%$ and goodness of fit (GoF) of 2.46. The fitted ordered magnetic moment of $0.52(6)\mu_B$ is in good agreement with the above values obtained from the graph in Fig. 5(b). Hence, we reach the same conclusion as from the analysis of the powder data that the magnetic space group for the magnetically ordered state of SrCuTe₂O₆ below T_{N1} is $P4_132$.

The obtained ordered moment of $0.52(6)\mu_B$ is about half of the expected value of $1\mu_B$ for $S = 1/2$ (48% reduction), suggesting that frustration in the J_1 and J_2 bonds potentially induces spin fluctuations and significantly reduces the ordered moment. The reduction of ordered moments has been observed in ordered frustrated systems, KFe₃(OH)₆(SO₂)₄ ($S = 3/2$) with 24% reduction [46], and Cs₂Cu₃SnF₁₂ ($S = 1/2$) with 32% reduction [47]. In comparison, for PbCuTe₂O₆, where the dominant J_2 forms the frustrated hyper-kagome lattice, spin fluctuations are so large that the Néel state is

totally suppressed and a quantum spin liquid possibly emerges at low temperatures [24–26]. From the DFT calculations, the exchange interactions in $\text{PbCuTe}_2\text{O}_6$ are close to one another, which could enhance the frustration whereas the intrachain interaction J_3 in $\text{SrCuTe}_2\text{O}_6$ is, respectively, one order and two orders of magnitude larger than J_1 and J_2 , which could place $\text{SrCuTe}_2\text{O}_6$ away from the quantum spin liquid state even though structurally it is almost identical to $\text{PbCuTe}_2\text{O}_6$. Nevertheless, even though $\text{SrCuTe}_2\text{O}_6$ magnetically orders at low temperatures, the residue effect of the frustrated bonds remains and evidences in the reduced ordered moment. It would be interesting to investigate this subtle effect of frustration in spin dynamics of this system.

IV. CONCLUSION

Magnetization measurements on single-crystal $\text{SrCuTe}_2\text{O}_6$ reveal highly isotropic magnetic susceptibility along the three inequivalent directions [100], [110], and [111] in this cubic system. The value of the leading exchange interaction (J_3/k_B) estimated using a quantum Monte Carlo simulation on the 1D spin-chain model is between 50.1 and 52.5 K. The order parameter measured by neutron scattering confirms that the system magnetically orders below $T_{N1} = 5.25(5)$ K. However, from our zero-field neutron diffraction measurements, we are unable to confirm the field-dependent second transition at lower T_{N2} , which was previously observed in magnetization and heat capacity measurements. Further in-field neutron scattering measurements are required to investigate this second transition. Based on the neutron diffraction data on the powder and single-crystal samples, the magnetic structure in the Shubnikov space group $P4_132$, where the Cu^{2+} $S = 1/2$ spins antiferromagnetically align along the direction perpendicular to the spin chain, is proposed with the ordered magnetic moment of $0.52(6)\mu_B$. This work suggests the dominance of the intrachain interaction J_3 over the frustrated J_1 and J_2 bonds, and sheds light on the difference in magnetic ground states between $\text{SrCuTe}_2\text{O}_6$ and $\text{PbCuTe}_2\text{O}_6$. The 48% reduction of the ordered moment in $\text{SrCuTe}_2\text{O}_6$ points to the residual effect of frustration, which could have nontrivial influence on spin dynamics in this magnetically ordered system.

Note added in proof. After submitting this manuscript, we became aware of similar work [48], which was published in a public archive. Magnetic susceptibility measured on a single-crystal sample reported in Ref. [48] is consistent with our results. Reference [48] also confirms our reported magnetic structure with the reduced ordered magnetic moment. However, their measured value of the ordered moment is slightly lower than that reported in this work.

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APPENDIX A: ATOMIC COORDINATES OF $\text{SrCuTe}_2\text{O}_6$

The refined fractional atomic coordinates of $\text{SrCuTe}_2\text{O}_6$ were shown in Table III. The refinement of the powder (crushed crystalline sample) and single crystal data was performed using FULLPROF [31] and SHELXL [32], respectively.

APPENDIX B: TABLE OF MAGNETIC IRREDUCIBLE REPRESENTATIONS OF $\text{SrCuTe}_2\text{O}_6$

The magnetic structure of $\text{SrCuTe}_2\text{O}_6$ was analyzed using the irreducible representation theory. The analysis based on the symmetry of the underlying crystal structure (space group $P4_132$) was carried out using BASIREPS [31] in the FULLPROF software package. Since the magnetic Bragg reflections were observed on top of the structural reflections as shown in the inset of Fig. 5(a), the magnetic propagation vector \vec{k} is equal to (0,0,0). For the Wyckoff position 12d of magnetic Cu^{2+} ions with a total of 12 spins in the unit cell as shown in Figs. 1(b) and 7(c), the decomposition of the irreducible representations

TABLE III. Refined values of fractional coordinates of $\text{SrCuTe}_2\text{O}_6$ from powder and single crystal x-ray diffraction measured at room temperature, and powder neutron diffraction measured at 10 K.

Atom	Site	x/a	y/a	z/a
Powder x-ray diffraction				
Te	24e	0.3380(1)	0.9187(1)	0.0588(1)
Sr(1)	8c	0.0545(2)	0.0545(2)	0.0545(2)
Sr(2)	4b	0.375	0.625	0.125
Cu	12d	0.4762(1)	0.875	0.2738(3)
O(1)	24e	0.6635(9)	1.1271(9)	0.1761(9)
O(2)	24e	0.4404(9)	1.0205(9)	1.2210(8)
O(3)	24e	0.2222(9)	0.9781(10)	0.1302(11)
$R_p = 0.0666$, $R_{wp} = 0.0849$, GoF = 2.1				
Single-crystal x-ray diffraction				
Te	24e	0.33827(3)	0.91872(3)	0.05938(3)
Sr(1)	8c	0.05469(5)	0.05469(5)	0.05469(5)
Sr(2)	4b	0.375	0.625	0.125
Cu	12d	0.47567(7)	0.875	0.27433(7)
O(1)	24e	0.6710(4)	1.1273(4)	0.1785(4)
O(2)	24e	0.4382(4)	1.0171(4)	1.2284(4)
O(3)	24e	0.2224(5)	0.9767(6)	0.1305(5)
$R_1 = 0.0247$, $wR_2 = 0.0601$, GoF = 1.042				
Powder neutron diffraction				
Te	24e	0.3379(1)	0.9192(1)	0.0589(1)
Sr(1)	8c	0.0536(1)	0.0536(1)	0.0536(1)
Sr(2)	4b	0.375	0.625	0.125
Cu	12d	0.4760(1)	0.875	0.2741(1)
O(1)	24e	0.6702(1)	1.1271(1)	0.1795(1)
O(2)	24e	0.4391(1)	1.0163(1)	1.2271(1)
O(3)	24e	0.2220(1)	0.9766(2)	0.1297(1)
$R_p = 0.0372$, $R_{wp} = 0.0487$, GoF = 2.5				

TABLE IV. Magnetic irreducible representations and their basis vectors for Cu1(x, y, z), Cu2($-x + 1/2, -y + 1, z + 1/2$), Cu3($-x + 1, y - 1/2, -z + 1/2$), Cu4($x + 1/2, -y + 3/2, -z + 1$), Cu5(z, x, y), Cu6($z + 1/2, -x + 1/2, -y + 1$), Cu7($-z + 1/2, -x + 1, y - 1/2$), Cu8($-z + 1, x + 1/2, -y + 3/2$), Cu9(y, z, x), Cu10($-y + 1, z + 1/2, -x + 1/2$), Cu11($y - 1/2, -z + 1/2, -x + 1$), Cu12($-y + 3/2, -z + 1, x + 1/2$).

IRs	BV	Cu1	Cu2	Cu3	Cu4	Cu5	Cu6	Cu7	Cu8	Cu9	Cu10	Cu11	Cu12
Γ_1	ψ_1	(10-1) ^a	(-10-1)	(-101)	(101)	(-110)	(-1-10)	(1-10)	(110)	(0-11)	(0-1-1)	(01-1)	(011)
Γ_2	ψ_1	(101)	(-101)	(-10-1)	(10-1)	(110)	(1-10)	(-1-10)	(-110)	(011)	(01-1)	(0-1-1)	(0-11)
	ψ_2	(010)	(0-10)	(010)	(0-10)	(001)	(00-1)	(001)	(00-1)	(100)	(-100)	(100)	(-100)
Γ_3^b	ψ_1	(100)	(-100)	(-100)	(100)	(0- $\frac{1}{2}$ 0)	(0 $\frac{1}{2}$ 0)	(0 $\frac{1}{2}$ 0)	(0- $\frac{1}{2}$ 0)	(00- $\frac{1}{2}$)	(00 $\frac{1}{2}$)	(00 $\frac{1}{2}$)	(00- $\frac{1}{2}$)
		(000)	(000)	(000)	(000)	(0- $\frac{\sqrt{3}}{2}$ 0)	(0 $\frac{\sqrt{3}}{2}$ 0)	(0 $\frac{\sqrt{3}}{2}$ 0)	(0- $\frac{\sqrt{3}}{2}$ 0)	(00 $\frac{\sqrt{3}}{2}$)	(00- $\frac{\sqrt{3}}{2}$)	(00- $\frac{\sqrt{3}}{2}$)	(00 $\frac{\sqrt{3}}{2}$)
	ψ_2	(010)	(0-10)	(010)	(0-10)	(00- $\frac{1}{2}$)	(00 $\frac{1}{2}$)	(00- $\frac{1}{2}$)	(00 $\frac{1}{2}$)	(- $\frac{1}{2}$ 00)	($\frac{1}{2}$ 00)	(- $\frac{1}{2}$ 00)	($\frac{1}{2}$ 00)
		(000)	(000)	(000)	(000)	(00- $\frac{\sqrt{3}}{2}$)	(00 $\frac{\sqrt{3}}{2}$)	(00- $\frac{\sqrt{3}}{2}$)	(00 $\frac{\sqrt{3}}{2}$)	($\frac{\sqrt{3}}{2}$ 00)	(- $\frac{\sqrt{3}}{2}$ 00)	($\frac{\sqrt{3}}{2}$ 00)	(- $\frac{\sqrt{3}}{2}$ 00)
	ψ_3	(001)	(001)	(00-1)	(00-1)	(- $\frac{1}{2}$ 00)	(- $\frac{1}{2}$ 00)	($\frac{1}{2}$ 00)	($\frac{1}{2}$ 00)	(0- $\frac{1}{2}$ 0)	(0- $\frac{1}{2}$ 0)	(0 $\frac{1}{2}$ 0)	(0 $\frac{1}{2}$ 0)
		(000)	(000)	(000)	(000)	(- $\frac{\sqrt{3}}{2}$ 00)	(- $\frac{\sqrt{3}}{2}$ 00)	($\frac{\sqrt{3}}{2}$ 00)	($\frac{\sqrt{3}}{2}$ 00)	(0 $\frac{\sqrt{3}}{2}$ 0)	(0 $\frac{\sqrt{3}}{2}$ 0)	(0- $\frac{\sqrt{3}}{2}$ 0)	(0- $\frac{\sqrt{3}}{2}$ 0)
	ψ_4	(00 $\frac{1}{2}$)	(00 $\frac{1}{2}$)	(00- $\frac{1}{2}$)	(00- $\frac{1}{2}$)	($\frac{1}{2}$ 00)	($\frac{1}{2}$ 00)	(- $\frac{1}{2}$ 00)	(- $\frac{1}{2}$ 00)	(0-10)	(0-10)	(010)	(010)
		(00- $\frac{\sqrt{3}}{2}$)	(00- $\frac{\sqrt{3}}{2}$)	(00 $\frac{\sqrt{3}}{2}$)	(00 $\frac{\sqrt{3}}{2}$)	($\frac{\sqrt{3}}{2}$ 00)	($\frac{\sqrt{3}}{2}$ 00)	(- $\frac{\sqrt{3}}{2}$ 00)	(- $\frac{\sqrt{3}}{2}$ 00)	(000)	(000)	(000)	(000)
	ψ_5	(0 $\frac{1}{2}$ 0)	(0- $\frac{1}{2}$ 0)	(0 $\frac{1}{2}$ 0)	(0- $\frac{1}{2}$ 0)	(00 $\frac{1}{2}$)	(00- $\frac{1}{2}$)	(00 $\frac{1}{2}$)	(00- $\frac{1}{2}$)	(-100)	(100)	(-100)	(100)
		(0- $\frac{\sqrt{3}}{2}$ 0)	(0 $\frac{\sqrt{3}}{2}$ 0)	(0- $\frac{\sqrt{3}}{2}$ 0)	(0 $\frac{\sqrt{3}}{2}$ 0)	(00 $\frac{\sqrt{3}}{2}$)	(00- $\frac{\sqrt{3}}{2}$)	(00 $\frac{\sqrt{3}}{2}$)	(00- $\frac{\sqrt{3}}{2}$)	(000)	(000)	(000)	(000)
	ψ_6	($\frac{1}{2}$ 00)	(- $\frac{1}{2}$ 00)	(- $\frac{1}{2}$ 00)	($\frac{1}{2}$ 00)	(0 $\frac{1}{2}$ 0)	(0- $\frac{1}{2}$ 0)	(0- $\frac{1}{2}$ 0)	(0 $\frac{1}{2}$ 0)	(00-1)	(001)	(001)	(00-1)
		(- $\frac{\sqrt{3}}{2}$ 00)	($\frac{\sqrt{3}}{2}$ 00)	($\frac{\sqrt{3}}{2}$ 00)	(- $\frac{\sqrt{3}}{2}$ 00)	(0 $\frac{\sqrt{3}}{2}$ 0)	(0- $\frac{\sqrt{3}}{2}$ 0)	(0- $\frac{\sqrt{3}}{2}$ 0)	(0 $\frac{\sqrt{3}}{2}$ 0)	(000)	(000)	(000)	(000)
Γ_4	ψ_1	(100)	(-100)	(100)	(-100)	(000)	(000)	(000)	(000)	(0-10)	(010)	(0-10)	(010)
	ψ_2	(010)	(0-10)	(0-10)	(010)	(000)	(000)	(000)	(000)	(-100)	(-100)	(100)	(100)
	ψ_3	(001)	(001)	(001)	(001)	(000)	(000)	(000)	(000)	(00-1)	(00-1)	(00-1)	(00-1)
	ψ_4	(000)	(000)	(000)	(000)	(-110)	(110)	(1-10)	(-1-10)	(000)	(000)	(000)	(000)
	ψ_5	(00-1)	(001)	(0-1)	(001)	(010)	(0-10)	(010)	(0-10)	(000)	(000)	(000)	(000)
	ψ_6	(0-10)	(0-10)	(010)	(010)	(001)	(00-1)	(00-1)	(001)	(000)	(000)	(000)	(000)
	ψ_7	(-100)	(-100)	(-100)	(-100)	(100)	(100)	(100)	(100)	(000)	(000)	(000)	(000)
	ψ_8	(000)	(000)	(000)	(000)	(000)	(000)	(000)	(000)	(0-11)	(011)	(01-1)	(0-1-1)
	ψ_9	(000)	(000)	(000)	(000)	(-100)	(100)	(-100)	(100)	(001)	(00-1)	(001)	(00-1)
	ψ_{10}	(000)	(000)	(000)	(000)	(00-1)	(00-1)	(001)	(001)	(100)	(-100)	(-100)	(100)
	ψ_{11}	(000)	(000)	(000)	(000)	(0-10)	(0-10)	(0-10)	(0-10)	(010)	(010)	(010)	(010)
	ψ_{12}	(10-1)	(101)	(-101)	(-10-1)	(000)	(000)	(000)	(000)	(000)	(000)	(000)	(000)
Γ_5	ψ_1	(100)	(-100)	(100)	(-100)	(000)	(000)	(000)	(000)	(010)	(0-10)	(010)	(0-10)
	ψ_2	(010)	(0-10)	(0-10)	(010)	(000)	(000)	(000)	(000)	(100)	(100)	(-100)	(-100)
	ψ_3	(001)	(001)	(001)	(001)	(000)	(000)	(000)	(000)	(001)	(001)	(001)	(001)
	ψ_4	(000)	(000)	(000)	(000)	(110)	(-110)	(-1-10)	(1-10)	(000)	(000)	(000)	(000)
	ψ_5	(000)	(000)	(000)	(000)	(001)	(001)	(001)	(001)	(000)	(000)	(000)	(000)
	ψ_6	(001)	(00-1)	(001)	(00-1)	(010)	(0-10)	(010)	(0-10)	(000)	(000)	(000)	(000)
	ψ_7	(010)	(010)	(0-10)	(0-10)	(001)	(00-1)	(00-1)	(001)	(000)	(000)	(000)	(000)
	ψ_8	(100)	(100)	(100)	(100)	(100)	(100)	(100)	(100)	(000)	(000)	(000)	(000)
	ψ_9	(000)	(000)	(000)	(000)	(000)	(000)	(000)	(000)	(011)	(0-11)	(0-1-1)	(01-1)
	ψ_{10}	(000)	(000)	(000)	(000)	(000)	(000)	(000)	(000)	(100)	(100)	(100)	(100)
	ψ_{11}	(000)	(000)	(000)	(000)	(100)	(-100)	(100)	(-100)	(001)	(00-1)	(001)	(00-1)
	ψ_{12}	(000)	(000)	(000)	(000)	(001)	(001)	(00-1)	(00-1)	(100)	(-100)	(-100)	(100)
	ψ_{13}	(000)	(000)	(000)	(000)	(010)	(010)	(010)	(010)	(010)	(010)	(010)	(010)
	ψ_{14}	(101)	(10-1)	(-10-1)	(-101)	(000)	(000)	(000)	(000)	(000)	(000)	(000)	(000)
	ψ_{15}	(010)	(010)	(010)	(010)	(000)	(000)	(000)	(000)	(000)	(000)	(000)	(000)

^aA parenthesis represents ($m_x m_y m_z$).

^bThe basis vectors for Γ_3 are complex with the first and second rows denoting the real and imaginary parts, respectively.

(IRs) can be described by

$$\Gamma = 1\Gamma_1^{(1)} + 2\Gamma_2^{(1)} + 3\Gamma_3^{(2)} + 4\Gamma_4^{(3)} + 5\Gamma_5^{(3)}, \quad (\text{B1})$$

where the basis vectors for Γ_1 , Γ_2 , Γ_3 , Γ_4 , and Γ_5 are given in Table IV. Γ_1 and Γ_2 are one dimensional with one and two basis vector(s), respectively. On the other hand, Γ_3 are of two

dimensions with six basis vectors whereas Γ_4 and Γ_5 are of three dimensions with 12 and 15 basis vectors, respectively. We assume that there is only one order parameter for the magnetic transition in SrCuTe₂O₆ and hence based on the Landau theory [49], the magnetic structure of the low-temperature phase corresponds to a single IR.

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