

Change of magnetic ground state by light electron doping in CeOs₂Al₁₀

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The effect of Ir substitution for Os in CeOs₂Al₁₀, with an unusually high Néel temperature of $T^* \sim 28.5$ K, has been studied by high-resolution neutron diffraction and magnetization measurements. A small amount of Ir ($\sim 8\%$) results in a pronounced change of the magnetic structure of the Ce sublattice. The induced magnetic ground state is controlled by single ion anisotropy and implies an antiferromagnetic arrangement of the Ce moments along the a axis, as expected from the anisotropy of the paramagnetic susceptibility. The value of the ordered moments, $0.92(1)\mu_B$, is substantially bigger than in the undoped compound, whereas the transition temperature is reduced down to 21 K. A comparison of the observed phenomena with the recently studied CeRu_{1.9}Rh_{0.1}Al₁₀ system, exhibiting similar behavior [Kondo *et al.*, *J. Phys. Soc. Jpn.* **82**, 054709 (2013)], strongly suggests the electron doping as the main origin of the ground state changes. This provides additional ways for exploring the anomalous magnetic properties of the Ce(Ru/Os)₂Al₁₀ compounds.

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The controversial interpretation of the phase transition recently found in the Kondo semiconductors CeRu₂Al₁₀ (Ref. 1) and CeOs₂Al₁₀ (Refs. 2 and 3) initiated their extensive study. After some debate,^{2–8} the nature of the transition was proven to be magnetic due to the ordering of the Ce sublattice.^{9–11} This ordering, however, is rather unusual and involves a very small ordered moment along a direction not expected from the anisotropy of the static susceptibility (χ), measured above the transition temperature. The anisotropy is very large and implies the a -crystallographic direction as the easy magnetic axis ($\chi_a > \chi_c > \chi_b$), whereas the neutron diffraction indicates the ordered moment, $0.3\text{--}0.4\mu_B$, to be along the c direction.^{9,12,13} In addition, the ordering temperature is unexpectedly high, $T^* \sim 28$ K, if one takes into account the moment value and the large distance (>5.2 Å) separating the Ce ions in the structure. Finally, a spin gap formation has been observed in inelastic neutron scattering experiments below the transition temperature in CeRu₂Al₁₀ and CeOs₂Al₁₀ as well as in the paramagnetic state of CeFe₂Al₁₀.^{7,11,14,15}

Aiming to elucidate the local electronic structure of Ce, Strigari *et al.*^{16,17} undertook polarization-dependent soft x-ray absorption measurements combined with magnetization data for both CeRu₂Al₁₀ and CeOs₂Al₁₀. The proposed crystal field ground state wave function provides quantitative agreement with the measured moment in the high magnetic field applied along the a direction and the small ordered moment along the c direction found in the neutron diffraction experiments. The proposed ground state has been confirmed directly from the inelastic neutron scattering measurements on CeOs₂Al₁₀, which reveals two crystal field excitations at 40 and 60 meV.¹⁸ The mechanisms providing the high ordering temperature and the gap formation are, however, unclear so far. Kimura *et al.*¹⁹ reported anisotropic changes in the electronic structure of CeOs₂Al₁₀ from polarized optical conductivity measurements and suggested that a charge density wave formation above T^* is the primary instability which then induces anomalous magnetic ordering at T^* .

Very recently, Kondo *et al.*²⁰ reported a remarkable change in the magnetic properties of CeRu₂Al₁₀ induced by a small

amount ($\sim 5\%$) of Rh substitution onto Ru. In particular, a spin-flop-like transition has been observed in a magnetic field ~ 13 T applied along the a direction, indicating a different ground state with the ordered moments being along this axis, as expected from the paramagnetic susceptibility. The observed effect has been attributed to the extra $4d$ electrons carried by Rh ions, suppressing the anisotropic character of the c - f hybridization and promoting a localized state for the $4f$ electrons of Ce.

This electron-doping induced ground state is of great interest and the primary questions immediately spring to mind: Does the doping result in any structural changes? Does it affect the magnetic propagation vector and the ordered moment size? In the present Rapid Communication, we address these matters for the CeOs_{1.84}Ir_{0.16}Al₁₀ system by means of high-resolution neutron diffraction. We found that the small amount of Ir carrying extra $5d$ electrons indeed results in the ground state change with a considerably bigger ordered moment in comparison with the undoped counterpart. The moments are found to be along the a axis, indicating the predominance of single ion anisotropy with an antiferromagnetic arrangement that is similar to that reported before for the CeOs₂Al₁₀ and CeRu₂Al₁₀ compounds.

The polycrystalline sample of CeOs_{1.84}Ir_{0.16}Al₁₀ was prepared by ultrahigh-purity argon arc melting of stoichiometric quantities of the starting elements. The sample was annealed at 850° C for 7 days for homogenization. A single crystal of the same composition was grown by the Al self-flux method.³ The neutron powder diffraction data were collected on the WISH time-of-flight diffractometer at the ISIS Facility of the Rutherford Appleton Laboratory, UK.²¹ The sample (6 g) was loaded into a cylindrical 6 mm vanadium can and measured on warming between 1.5 and 30 K in 3 K steps using an Oxford Instrument cryostat. The crystal and magnetic structure Rietveld refinements were performed using the FULLPROF program²² against the data measured in detector banks at average 2θ values of 58°, 90°, 122°, and 154°, each covering 32° of the scattering plane. Magnetic susceptibility was measured using a superconducting quantum interference

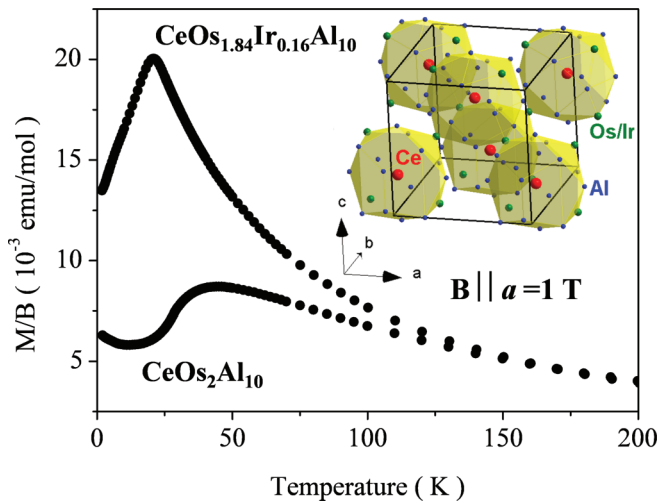


FIG. 1. (Color online) Magnetic susceptibility of $\text{Ce}(\text{Os}_{1-x}\text{Ir}_x)_2\text{Al}_{10}$ compositions, measured along the a direction, as a function of temperature. The inset shows a schematic representation of the crystal structure of these compositions as a stacking of the $\text{CeAl}_{16}(\text{Os/Ir})_4$ polyhedron cages.

device (SQUID) magnetometer (Quantum Design MPMS) on both $\text{CeOs}_2\text{Al}_{10}$ and $\text{CeOs}_{1.84}\text{Ir}_{0.16}\text{Al}_{10}$ single crystals.

The magnetic data obtained for the lightly Ir-doped crystal revealed a pronounced change in the low temperature susceptibility measured along the a direction (Fig. 1). The changes are related to both the paramagnetic and the ordered phases, indicating behavior similar to that recently observed by Kondo *et al.*²⁰ in the $\text{CeRu}_{1.9}\text{Rh}_{0.1}\text{Al}_{10}$ system. Above ~ 150 K, the paramagnetic susceptibilities for both doped and undoped crystals are very close. Below this temperature, the susceptibility for the undoped crystal starts deviating from the Curie-Weiss behavior, whereas the Ir-doped crystal keeps it down to lower temperatures. Around $T^* \sim 21$ K, the susceptibility passes through a maximum, indicating the phase transition. The critical temperature is essentially lower than in the undoped compound and at present it is unknown how it varies with the Ir content.

To directly explore the magnetic ground state of the electron doped $\text{CeOs}_{1.84}\text{Ir}_{0.16}\text{Al}_{10}$ system, a high-resolution neutron diffraction study has been performed. Although the doping level is small, possible structural changes were explored in the refinement procedure. The parent Ir-free compound crystallizes into the orthorhombic $\text{YbFe}_2\text{Al}_{10}$ structure type with the $Cmcm$ space group.^{23–25} The structure encapsulates Ce into polyhedron cages formed by sixteen Al and four Os atoms (Fig. 1 inset), providing a large nearest-neighbor Ce-Ce distance. This structural model works well for all patterns collected at temperatures below and above the transition [Fig. 2(c)], so within the resolution limit of our diffraction experiment no structural modifications either induced compositionally or temperature induced could be observed. A small amount of the $\text{Os}_4\text{Al}_{13}$ impurity phase has been identified and included into the refinement. The structural parameters obtained at $T = 30$ K are summarized in Table I. Evaluation of the unit cell parameters and the Ce-Ce bond distances also did not reveal any notable anomalies across the magnetic transition, indicating a weak magnetoelastic coupling in the system.

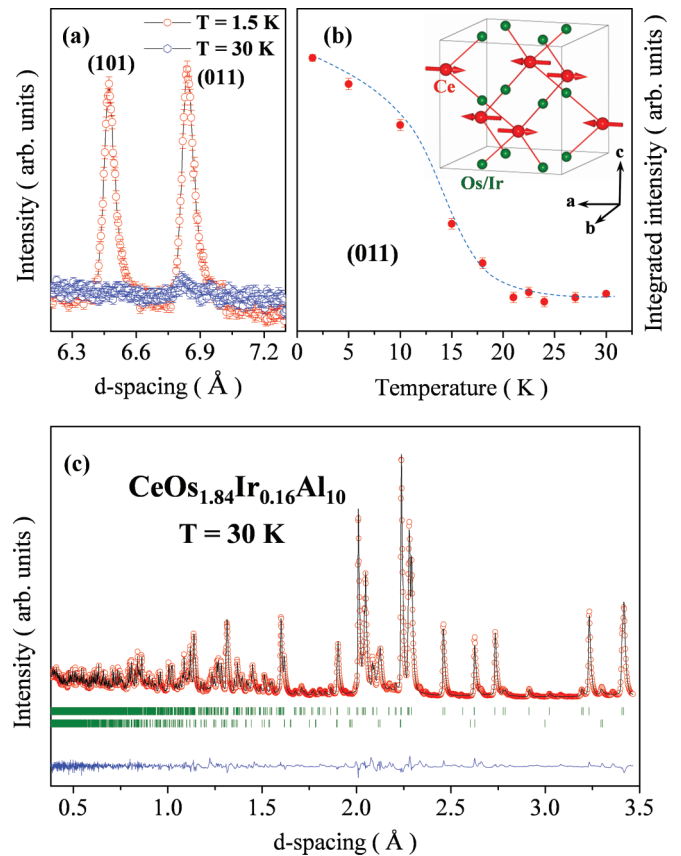


FIG. 2. (Color online) (a) A part of the neutron diffraction patterns near the (101) and (011) magnetic reflections collected at different temperatures. (b) Integrated intensity of the (011) magnetic peak as a function of temperature. The dotted line is a guide to the eyes. The inset shows the magnetic ground state of $\text{CeOs}_{1.84}\text{Ir}_{0.16}\text{Al}_{10}$ (for clarity, only Ce and Os/Ir atoms are shown). (c) Rietveld refinement of the neutron powder diffraction pattern collected at the backscattering detector bank (average scattering angle is 154°) of the WISH diffractometer. The circle symbols (red) and solid line represent the experimental and calculated intensities, respectively, and the line below (blue) is the difference between them. Tick marks indicate the positions of Bragg peaks for the $\text{CeOs}_{1.84}\text{Ir}_{0.16}\text{Al}_{10}$ (top) and $\text{Os}_4\text{Al}_{13}$ (bottom) phases.

Magnetic Bragg scattering appears below $T^* \sim 21$ K, resulting in a set of additional reflections clearly visible in a large d -spacing region [Fig. 2(a)]. Their appearance correlates with a background reduction at the low momentum transfer part of the patterns. This behavior is expected in the case of a long-range magnetic ordering due to decreasing the paramagnetic scattering. All the magnetic peaks can be indexed with the $\mathbf{k} = (1,0,0)$ propagation vector, identical to the undoped $\text{CeOs}_2\text{Al}_{10}$ compound. The integrated intensities of these peaks as a function of temperature demonstrate a mean-field-like behavior consistent with the continuous nature of the transition. To refine the magnetic structure, a difference between patterns collected at 1.5 and 30 K was taken to better evaluate the fitting quality introduced by different models. The refinement procedure was conducted by the symmetry-based approach, implying classification of the symmetrized pseudovector modes, localized on the Ce

TABLE I. Structural parameters of $\text{CeOs}_{1.84}\text{Ir}_{0.16}\text{Al}_{10}$ refined from the neutron diffraction data collected at $T = 30$ K in the orthorhombic $Cmcm$ [$a = 9.1193(2)$ Å, $b = 10.2554(2)$ Å, $c = 9.1657(2)$ Å, $R_{\text{Bragg}} = 4.18\%$] space group. Occupancies for all the atoms in the refinement procedure were fixed to the nominal chemical content.

Atom	Site	x	y	z	B_{iso}
Ce	4c	0	0.1260(6)	0.25	0.23(8)
Os/Ir	8d	0.25	0.25	0	0.40(4)
Al1	8g	0.2216(7)	0.3661(7)	0.25	0.2(1)
Al2	8g	0.3512(6)	0.1305(7)	0.25	0.4(1)
Al3	8f	0	0.1587(7)	0.6005(6)	0.2(1)
Al4	8f	0	0.3764(7)	0.0485(6)	0.2(1)
Al5	8e	0.2231(6)	0	0	0.3(1)

and Os/Ir positions, accordingly irreducible representations of the parent $Cmcm$ space group.^{26,27} A unique solution involving the single Y_2^- representation (in the Miller and Love notations²⁸) and nonzero moments only for the Ce sites [Fig. 2(b) inset] has been identified based on the quality fitting of the magnetic intensities [Fig. 3(a)]. The obtained model has the C_{pmcm}' magnetic symmetry (No. 63.13.523 in Litvin's classification scheme²⁹) and is different from those reported for undoped $\text{CeOs}_2\text{Al}_{10}$ and $\text{CeRu}_2\text{Al}_{10}$ by the direction of the magnetic moments, which are along the a axis in this ground state. The type of antiferromagnetic coupling violating the C -centering translation remains unchanged. In particular, the ferromagnetic arrangement of Ce moment in the zigzag chain along the c axis does not change between the pure and doped systems. This fact may suggest the importance of a magnetic interaction along the zigzag chain via the hybridization with Os in the middle. The refined value of the Ce magnetic moments is substantially bigger, $0.92(1)\mu_B$ versus $0.29(1)\mu_B$ in the undoped compound.¹² Refinement in the model with the moments aligned along the c axis yields a much worse agreement with the experimental data [Fig. 3(b)] and can be unambiguously ruled out.

Thus, the obtained result directly reveals the change of the ground state of the system upon Ir substitution. The fact that this remarkable change has been observed in both $\text{Ce}(\text{RuRh})_2\text{Al}_{10}$ and $\text{Ce}(\text{OsIr})_2\text{Al}_{10}$ series provides a strong evidence that its main driving force is the electron doping rather than some crystal chemical effects such as disorder, compositionally induced symmetry change, or chemical pressure. The value of the ordered moment, $0.92(1)\mu_B$, is in a good agreement with the high field magnetization measurements reported by Kondo *et al.*³⁰ in the undoped $\text{CeOs}_2\text{Al}_{10}$. The latter was found to be $0.95\mu_B$ along the easy a axis. The crystal field ground state wave function deduced by Strigari *et al.*¹⁷ yields a somewhat bigger moment along this axis, $\sim 1.35\mu_B$, indicating the presence of the Kondo screening effect. Finally, it should be pointed out that in spite of the bigger ordered moment and lower transition temperature in $\text{CeOs}_{1.84}\text{Ir}_{0.16}\text{Al}_{10}$, this composition still does not obey the de Gennes's scaling applied to the $\text{Ln}(\text{Os/Ru})_2\text{Al}_{10}$ series (Ln = lanthanide), revealing an anomalous mechanism providing the magnetic ordering in this class of Ce-based materials. A

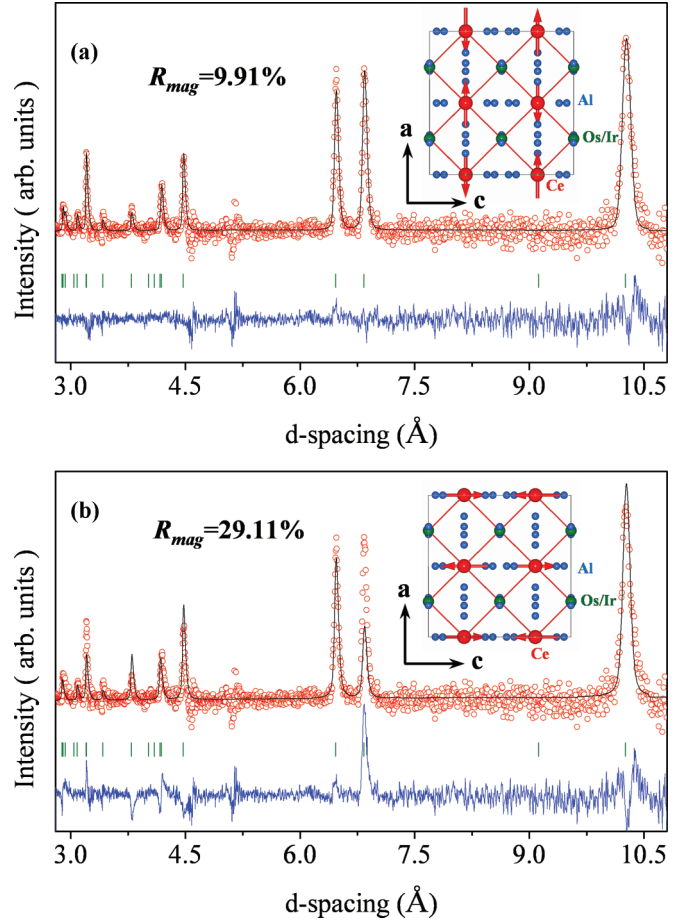


FIG. 3. (Color online) Rietveld refinements of the magnetic intensity of the $\text{CeOs}_{1.84}\text{Ir}_{0.16}\text{Al}_{10}$ composition obtained as a difference between the diffraction patterns collected at 1.5 and 30 K. The circle symbols (red) and solid line represent the experimental and calculated intensities, respectively, and the line below (blue) is the difference between them. Tick marks indicate the positions of Bragg peaks for the magnetic scattering with the ($\mathbf{k} = 1, 0, 0$) propagation vector. The refinement quality is demonstrated for two models shown as insets, with moments along the (a) a axis and (b) c axis.

comprehensive study of the electron-doping induced ground state can furnish useful insight into this mechanism. In particular, the presence or absence of the spin gap formation in the doped compounds is the primary point which has to be addressed. In addition, the more localized nature of the induced ground state can promote a magnetic order in the $\text{Ce}(\text{Fe}_{1-x}\text{Co}_x)_2\text{Al}_{12}$ system and provide better conditions for the experimental study of the crystal field effects using the direct inelastic neutron scattering technique.

In conclusion, a small substitution of Os with Ir in the $\text{CeOs}_2\text{Al}_{10}$ compound induces a pronounced change of the magnetic ground state. This magnetic ordering involves the antiferromagnetic arrangement of Ce moments along the a direction, as expected from the anisotropy of the paramagnetic susceptibility, with the ordered value being $0.92(1)\mu_B$. The transition temperature is sensitive to the doping as well and is $T^* \sim 21$ K for 8% of the Ir content. The main origin of the observed doping effect is the extra $5d$ electrons carried by Ir. The obtained result demonstrates a great sensitivity of the

Ce(Os/Ru)₂Al₁₀ systems to the carrier doping and provides additional ways to study their anomalous magnetic properties.

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