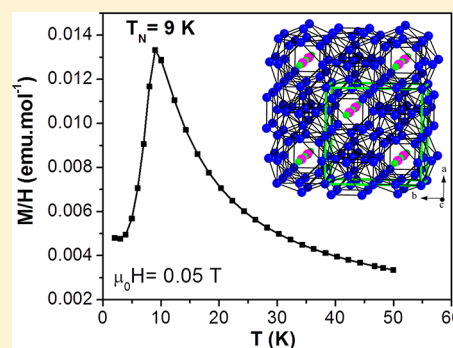


The Mg-Rich Phase NdNiMg<sub>15</sub>: Structural and Magnetic PropertiesEliane Al Asmar,<sup>†,‡,§</sup> Sophie Tencé,<sup>†,‡</sup> Jean-Louis Bobet,<sup>†,‡</sup> Bassem Ourane,<sup>†,‡</sup> Michel Nakhl,<sup>§</sup> Mirvat Zakhour,<sup>§</sup> and Etienne Gaudin<sup>\*,†,‡</sup><sup>†</sup>CNRS, ICMCB, UMR 5026, F-33600 Pessac, France<sup>‡</sup>Univ. Bordeaux, ICMCB, UMR 5026, F-33600 Pessac, France<sup>§</sup>LCPM/PR2N (EDST), Université Libanaise, Faculté des Sciences II, 90656, Jdeidet El Metn, Lebanon

**ABSTRACT:** The intermetallic NdNiMg<sub>15</sub> is the Mg-richest phase (more than 88 atom % of Mg) discovered in the Mg–Nd–Ni system. Its structure was determined by X-ray diffraction on single crystal with the following crystal data: tetragonal system, *P4/nmm*, *Z* = 2, *a* = 10.0602(1) Å, *c* = 7.7612(2) Å, *d*<sub>calc</sub> = 2.40 g·cm<sup>−3</sup>. Its structure is made of a three-dimensional framework of magnesium atoms showing channels filled by one-dimensional chain consisting of alternating Nd and Ni atoms along the *c*-axis. Anti-ferromagnetic ordering was observed with *T*<sub>N</sub> = 9 K, which is remarkably high considering the long distances between magnetic atoms, that is, Nd atoms. The effective magnetic moment  $\mu_{\text{eff}}$  is equal to 3.58  $\mu_{\text{B}}$ , which is consistent with magnetic Nd<sup>3+</sup> ions and weakly or nonmagnetic Ni atoms. Below *T*<sub>N</sub>, the *M*(*H*) curves show field-induced metamagnetic transitions at critical fields increasing with decreasing temperatures. The magnetic structure of NdNiMg<sub>15</sub> was determined from neutron powder diffraction data by considering the propagation vector *k* = (1/2 1/2 0). This magnetic structure consists in ferromagnetic chains along the *c*-axis of Nd atoms carrying moments, only separated by Ni atoms. The chains are ferromagnetically coupled within planes perpendicular to the [110] direction, and these planes are anti-ferromagnetically coupled to neighboring planes forming a checkerboard-like magnetic structure.



## INTRODUCTION

The interest in magnesium alloys has increased over the past 30 years, especially because magnesium is abundant in nature, cheap, lightweight, and has very diverse applications (automobile, aerospace, hydrogen storage, etc.). Numerous studies have been devoted to the search for new Mg-based light structural materials, mainly for automobile and aeronautic industry components<sup>1,2</sup> and for biomedical applications.<sup>3,4</sup> Besides these structural applications, magnesium and magnesium alloys are also good candidates for solid storage of hydrogen through the formation of metal hydrides.<sup>5,6</sup> Magnesium exhibits very high hydrogen storage capacity (i.e., 7.6 wt %, 110 gH/L by formation of MgH<sub>2</sub>), but its use remains limited due to the high operating temperature (~300 °C) and the heat released by the highly exothermic absorption reaction.<sup>5</sup> In addition, a new interest in the hydrolysis reaction between Mg or Mg alloys and water for hydrogen generation has recently emerged. This would be an interesting way to valorize Mg-based alloys wastes.<sup>7,8</sup> For all these reasons, a reinvestigation of the Mg-rich zone of many ternary (and quaternary) systems has been performed in the last years.

In the case of structural materials, extensive studies of the RE–Zn–Mg (RE = rare-earth) systems close to the Mg corner have led to the discovery of a new series of so-called long period stacking ordered (LPSO) phases.<sup>9</sup> The strength and ductility of Y–Zn–Mg alloys are strongly enhanced when they are produced by rapid solidification processing.<sup>10</sup> After this

process, intermetallic particles, so-called LPSO, precipitate and are distributed along the Mg grains boundaries. The LPSO phases are long period stacking variants of the structure of hexagonal close-packed (hcp)-Mg and the most frequently observed polytypes crystallizing with a 18R, 14H, or 10H structure.<sup>11,12</sup> The RE and Zn atoms are located in some of the layers, and they form Zn<sub>6</sub>RE<sub>8</sub> clusters. The Mg-rich corner of many RE–TM–Mg ternary systems (TM = transition metal) have been also studied for solid hydrogen storage, to combine the low weight of magnesium and the catalytic effect of rare-earth and transition metal.<sup>13–15</sup> Among these systems one can cite the La–Ni–Mg<sup>15,16</sup> and Nd–Ni–Mg<sup>17–21</sup> systems. The different studies show a decomposition of the intermetallic compounds into MgH<sub>2</sub>, binary rare-earth hydrides (e.g., LaH<sub>3</sub>), and ternary magnesium transition-metal hydrides (e.g., Mg<sub>2</sub>NiH<sub>4</sub>). An increase of the kinetics of the hydrogen sorption properties of Mg is observed, but no real breakthrough concerning the solid storage of hydrogen was achieved. The stabilization of these ternary phases upon hydrogenation is still an interesting challenge to overcome to find new candidates for solid hydrogen storage.

The reinvestigation of many ternary RE–TM–Mg systems has led to the discovery of numerous new very-rich Mg phases and has extended our knowledge of these ternary systems.

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Table 1. Atomic Positions and Equivalent Displacement Parameters of NdNiMg<sub>15</sub><sup>a</sup>

position	Wyck	<i>x</i>	<i>y</i>	<i>z</i>	<i>U</i> <sub>eq</sub> (Å <sup>2</sup> )
Nd1	2c	1/4	1/4	0.933 40(2)	0.009 68(4)
Ni2	2c	1/4	1/4	0.430 49(5)	0.014 73(8)
Mg3	4d	0	0	0	0.0166(2)
Mg4	8i	3/4	0.401 82(5)	0.150 49(7)	0.0155(1)
Mg5	2b	3/4	1/4	1/2	0.0317(3)
Mg6	8j	0.910 21(4)	0.910 21(4)	0.390 55(7)	0.0172(1)
Mg7	8i	3/4	0.980 25(5)	0.732 97(7)	0.0162(1)

<sup>a</sup>*P4*/*mmm* space group (origin choice 2) *a* = 10.0602(1) and *c* = 7.7612(2) Å.

Some of these ternary phases are disordered<sup>22,11</sup> or derive from binary phases.<sup>23,24</sup> For instance, Gd<sub>13</sub>Ni<sub>9</sub>Mg<sub>78</sub><sup>22</sup> crystallizes with a modulated structure that derives from an average cubic structure, and LaCuMg<sub>8</sub><sup>23</sup> crystallizes in the La<sub>2</sub>Mg<sub>17</sub>-type structure with Cu atoms located in interstitial sites or partially substituted to Mg atoms on one atomic position. Among the ordered ones one can cite NdNiMg<sub>5</sub>,<sup>18</sup> Mg<sub>30</sub>Co<sub>2</sub>Y<sub>9</sub>,<sup>25</sup> Ce<sub>2</sub>Ru<sub>4</sub>Mg<sub>17</sub>,<sup>26</sup> Nd<sub>16</sub>Mg<sub>96</sub>Ni<sub>12</sub>,<sup>21</sup> Nd<sub>4</sub>Ni<sub>8</sub>Mg<sub>80</sub>,<sup>20</sup> and the title compound NdNiMg<sub>15</sub><sup>27</sup> with increased molar ratio of magnesium (from 71.4% to 88.2%). Except NdNiMg<sub>5</sub>, which crystallizes with an ordered variant of the Nd<sub>3</sub>Co<sub>3</sub>Ga-type structure, all these compounds crystallize with new structural types. Only the magnetic properties of Ce<sub>2</sub>Ru<sub>4</sub>Mg<sub>17</sub> and NdNiMg<sub>5</sub> have been studied. No magnetic ordering was observed down to 3 K for the former compound, and this has been correlated to an almost tetravalent state of Ce atoms. An anti-ferromagnetic ordering is observed below *T*<sub>N</sub> = 12 K for NdNiMg<sub>5</sub>, and the value of the effective magnetic moment confirms that the prevalent contribution to magnetism comes from the Nd atoms. Among the LPSO phases discovered, recently Kishida et al. have determined the crystal structure of three phases with composition Y<sub>4</sub>Zn<sub>3</sub>Mg<sub>23</sub>, Y<sub>4</sub>Zn<sub>3</sub>Mg<sub>29</sub>, and Y<sub>4</sub>Zn<sub>3</sub>Mg<sub>35</sub>.<sup>12</sup> Currently, NdNiMg<sub>15</sub> is the richest Mg-ordered ternary crystalline phase ever discovered. In a recent paper we presented the interest of this compound for hydrogen generation.<sup>27</sup> The crystal structure, magnetic properties, and structure of this peculiar compound are reported herein.

## EXPERIMENTAL SECTION

The starting materials for the preparation of the sample were neodymium pieces (ZLX Tech, >99.9%), a nickel rod (Strem Chemicals, >99.9%), and a magnesium rod (Alpha Aesar, >99.8%). To avoid oxide impurities, the surfaces of the magnesium and neodymium pieces were carefully cleaned with abrasive tools. The Nd, Ni, and Mg elements were then weighed in the 1:1:15 atomic ratio and sealed in small tantalum ampules under an argon pressure of ca. 800 mbar. The tantalum ampule was purified before with magnesium sponge (673 K). Then, the ampules were placed in a high-frequency furnace under argon and heated at ~1773 K and kept at that temperature for 2 min. The tantalum ampule was then enclosed in an evacuated quartz tube and annealed at 973 K for 10 days. The temperature was decreased to room temperature at 6 K/h.

A crystal suitable for single-crystal X-ray diffraction was selected on the basis of the size and sharpness of the diffraction spots. Data collection was performed using a Bruker Kappa Apex II diffractometer (Mo radiation) at room temperature. Numerical absorption correction using the shape of the crystal (face indexed with the help of the video microscope) was made with SADABS-2014/5<sup>28</sup> software. The unit cell is tetragonal with lattice parameters *a* = 10.0602(1) and *c* = 7.7612(2) Å. The space group and structural determinations were performed with the Jana2006 program package.<sup>29</sup> The extinction conditions observed agree with the *P4*/*mmm* space group (the origin choice 2 was chosen for comparison with the

isotype structures). All the atomic positions were located using the SUPERFLIP program.<sup>30</sup> Refinement of the occupancy parameters showed full occupancy within two standard deviations. At the end of the refinement, with anisotropic atomic displacement parameters for all atomic positions, the reliability factors were *R*[*F*<sup>2</sup> > 2σ(*F*<sup>2</sup>)] = 0.012 and *wR*(*F*<sub>2</sub>) = 0.030, *S* = 1.00 for 31 parameters, 805 independent reflections, and with residual electron density in the range [−0.33, +0.51 e·Å<sup>−3</sup>]. Details of data collections and structure refinement can be found in the crystallographic information file (CIF) deposited in the Fachinformationszentrum Karlsruhe (FIZ, 76344 Eggenstein-Leopoldshafen Germany); with reference CSD 434727. The atomic positions are given in Table 1.

Magnetization measurements were performed using a superconducting quantum interference device (SQUID) magnetometer from Quantum Design for temperatures from 1.8 to 360 K and magnetic fields up to 7 T.

Neutron diffraction experiments were performed at the Laue Langevin Institute (ILL) on the D1B diffractometer with the use of a wavelength of 2.524 Å. A cylindrical vanadium sample holder of 6 mm inner diameter was used to collect neutron data. The FullProf Suite program was used for nuclear and magnetic structure refinements using the Rietveld method.<sup>31</sup>

## RESULTS AND DISCUSSION

**Crystal Structure.** The structure of NdNiMg<sub>15</sub> is made of a three-dimensional (3D) network of magnesium atoms giving rise to channels along the *c* axis, where chains made of nickel and neodymium atoms are located (Figure 1). The channels consist in infinite face-sharing chains of Mg<sub>8</sub> square antiprism filled alternatively by Ni and Nd atoms, and a connection between the chains of Mg<sub>8</sub> square prisms through Mg6–Mg6 bonds can be considered (Figure 2a). The Ni–Nd distances observed in the chain, equal to 3.858 and 3.903 Å (Table 2), are very long compared to the sum of the metallic radii, 3.07 Å

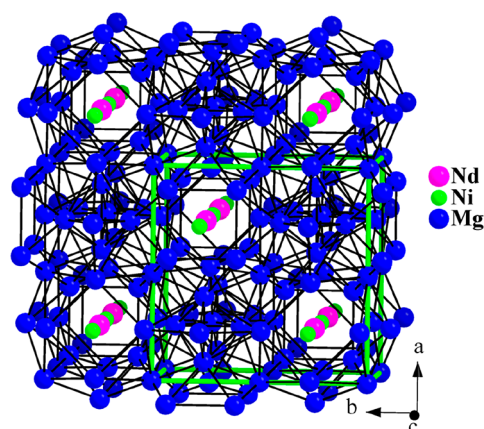
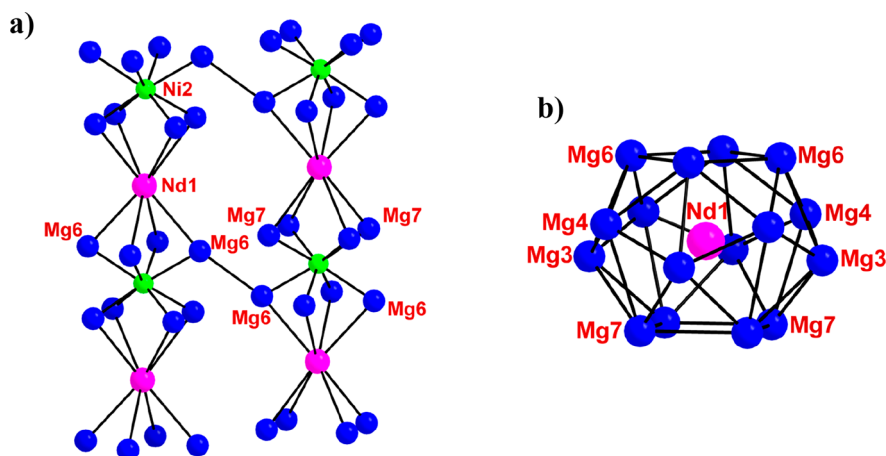


Figure 1. Perspective view of the structure of NdNiMg<sub>15</sub>. The Mg–Mg bonds are drawn (Mg–Mg distance lower than 3.3 Å).



**Figure 2.** (a) Nd–Ni chains of atoms with Nd–Mg, Ni–Mg, and Mg6–Mg6 bonds are drawn. (b) Environment of Nd1 position by considering distances Nd–Mg lower than 3.6 Å.

**Table 2. Interatomic Distances (in Å) for NdNiMg<sub>15</sub>**

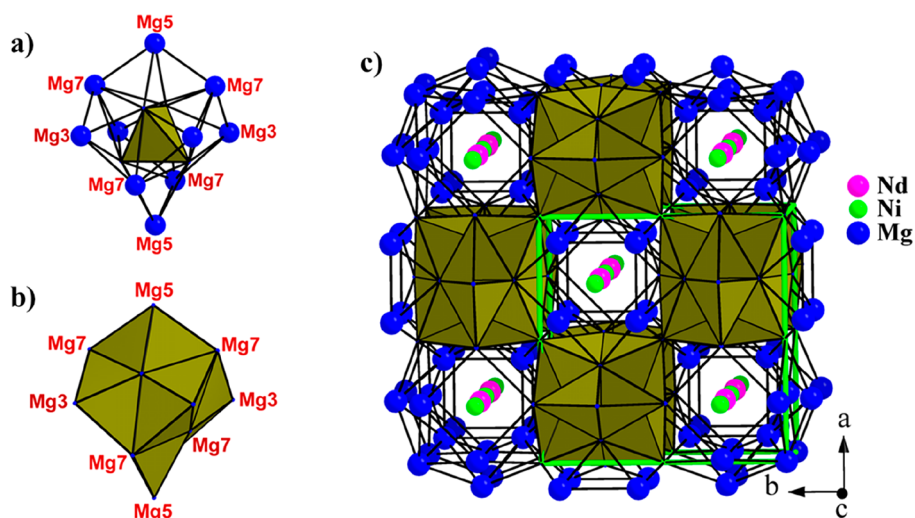
Nd1	Mg6 ×4	3.3936(5)	Mg5	Mg4 ×4	3.1130(5)
	Mg7 ×4	3.4743(5)		Mg7 ×4	3.2610(5)
	Mg4 ×4	3.5627(5)	Mg6	Ni2 ×1	2.6692(5)
	Mg3 ×4	3.5942(1)		Mg6 ×1	3.0682(6)
	Ni2 ×1	3.8580(4)		Mg4 ×2	3.1056(6)
Ni2 ×1	3.9032(5)	Mg7 ×2		3.1870(7)	
Ni2	Mg7 ×4	2.6410(6)	Mg6	Mg6 ×2	3.2235(6)
	Mg6 ×4	2.6692(5)		Mg3 ×1	3.2893(5)
	Nd1 ×1	3.8580(4)		Nd1 ×1	3.3936(5)
	Nd1 ×1	3.9032(5)			
Mg3			Mg7	Ni2 ×1	2.6410(6)
	Mg4 ×4	2.9437(3)		Mg6 ×2	3.1870(7)
	Mg7 ×4	3.2650(3)		Mg4 ×2	3.2427(5)
	Mg6 ×2	3.2893(5)		Mg5 ×1	3.2610(5)
Mg4				Mg3 ×2	3.2650(3)
	Mg3 ×2	2.9437(3)		Mg7 ×2	3.2758(5)
	Mg4 ×1	3.0547(8)		Mg4 ×1	3.4509(7)
	Mg6 ×2	3.1056(6)		Nd1 ×1	3.4743(5)
	Mg5 ×1	3.1130(5)			
	Mg4 ×2	3.1816(7)			
	Mg7 ×2	3.2427(5)			
	Mg7 ×1	3.4509(7)			

( $r_{\text{Nd}} = 1.82$  and  $r_{\text{Ni}} = 1.25$  Å<sup>32</sup>) and of the covalent radii, 2.79 Å ( $r_{\text{Nd}} = 1.64$  and  $r_{\text{Ni}} = 1.15$  Å<sup>33</sup>). Then, the bonding between Nd1 and Ni1 is expected to be rather weak. The Ni–Mg distances in the Mg<sub>8</sub> square antiprism, equal to 2.641 and 2.669 Å, are very close to the sum of the covalent radius of Ni and atomic radius of Mg, 2.75 Å ( $r_{\text{Mg}} = 1.60$  Å<sup>33</sup>). For the discussion of distances with Mg atoms, the covalent radius ( $r_{\text{Mg}} = 1.36$  Å<sup>33</sup>) has not been taken into account because of its too low value, which is not realistic in the case of intermetallic chemistry. Then the atomic radius of 1.60 Å has been considered, and moreover it corresponds also to the metallic radius.<sup>34</sup> By considering the two shortest Nd–Mg distances of 3.394 and 3.474 Å (Table 2), the Nd atom has eight Mg neighbors, which form a square antiprism. These distances are very close to the sum of the metallic radii ( $1.82 + 1.60 = 3.42$  Å). By slightly increasing the Nd–Mg distances up to 3.6 Å the coordination polyhedra becomes more complex (Figure 2b) with 16 Mg neighbors, 24 triangular faces, and 2 square faces.

In NdNiMg<sub>5</sub><sup>18</sup> different coordination polyhedra around Ni and Nd atoms are observed, but the Ni–Mg and Nd–Mg distances are very close to the ones observed for NdNiMg<sub>15</sub>. They range from 2.668 to 2.804 Å for Ni–Mg and from 3.292 to 3.416 Å for Nd–Mg (only the distances below 3.60 Å are considered). To describe the 3D network of magnesium atoms one can consider three roughly regular Mg<sub>4</sub> tetrahedra sharing faces or edges (Figure 3a,b). The most regular tetrahedron is [Mg<sub>4</sub>]<sub>4</sub> with Mg4–Mg4 distances equal to 3.055 and 3.182 Å and the angles between the center of the tetrahedra and the vertices equal to 105.2 and 111.7°. The two other types of tetrahedra are more distorted with in each case a long edge with a length equal to 3.265 Å for Mg3–Mg7 and 3.451 Å for Mg4–Mg7. A view of the structure with a drawing of these tetrahedra is shown in Figure 3c. The average distance between the center of these three tetrahedra and their vertices is equal to 1.94 Å, which is the same value as the one observed in Mg metal.<sup>34</sup> Moreover the Mg–Mg distances are ranging from 2.944 to 3.451 (Table 2) and are consistent with the value of 3.2 Å observed in Mg metal. All these structural features suggest the metallic character of the Mg network. The structure of NdNiMg<sub>15</sub> can be considered as an ordered variant of the V<sub>15</sub>Sb<sub>18</sub><sup>35</sup> structure type. Only a few compounds crystallize with this structural type; one can cite Ti<sub>8</sub>Bi<sub>9</sub>,<sup>36</sup> Zr<sub>2</sub>V<sub>6</sub>Sb<sub>9</sub>,<sup>37</sup> Ti<sub>5.42</sub>Mo<sub>2.58</sub>Sb<sub>9</sub>,<sup>38</sup> or Ce<sub>2</sub>Pd<sub>14</sub>Si.<sup>39</sup> It may be noticed that the ICSD database<sup>40</sup> has chosen the Zr<sub>2</sub>V<sub>6</sub>Sb<sub>9</sub> structure as structure type, even if the V<sub>15</sub>Sb<sub>18</sub> structure (or Bi<sub>8</sub>Ti<sub>9</sub>) has been published earlier. To our knowledge NdNiMg<sub>15</sub> is the first Mg-based phase crystallizing with this structural type.

The structures of the four very Mg-rich phases recently discovered in the Nd–Ni–Mg system can be described as made of a 3D metallic framework of magnesium atoms and of a covalent network of Nd–Ni atoms with lower dimensionality. A two-dimensional (2D) network is observed in NdNiMg<sub>5</sub>,<sup>18</sup> the Nd and Ni atoms forming ordered graphitic-type layer perpendicular to the *c*-axis and a one-dimensional (1D) network in NdNiMg<sub>15</sub> with chains alternatively made of Nd and Ni atoms running along the *c* axis. In the case of Nd<sub>4</sub>Ni<sub>8</sub>Mg<sub>80</sub><sup>20</sup> only Nd and Ni isolated atoms are observed, whereas in Nd<sub>16</sub>Ni<sub>12</sub>Mg<sub>96</sub><sup>21</sup> Nd<sub>4</sub>Ni<sub>2</sub> clusters and isolated Ni atoms are coexisting. As briefly discussed in the Introduction, the coexistence of a metallic Mg matrix and clusters of rare-earth and transition metal was previously observed in Mg-rich

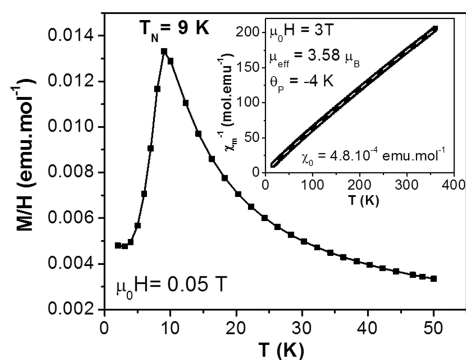




**Figure 3.** (a) Surrounding of a  $\text{Mg}_4$  tetrahedra drawn in green. (b) Same view as (a) with all  $\text{Mg}_4$  drawn in green. (c) View of the structure of  $\text{NdNiMg}_{15}$  with all  $\text{Mg}_4$  tetrahedra drawn in green.

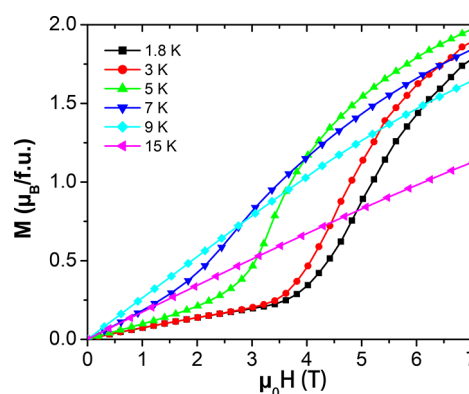
LPSO phases (see  $\text{Y-Zn-Mg}^{12}$  system, for instance) and led to drastic changes in mechanical properties.<sup>10</sup>

**Magnetic Properties.** The plot of the magnetization divided by the applied magnetic field  $M/H$  of  $\text{NdNiMg}_{15}$  is reported in Figure 4. On this curve, we clearly distinguish a



**Figure 4.** Temperature dependence of the magnetization divided by the magnetic field of  $\text{NdNiMg}_{15}$ . (inset) Inverse of the magnetic susceptibility vs temperature with the modified Curie–Weiss fit represented in dashed line.

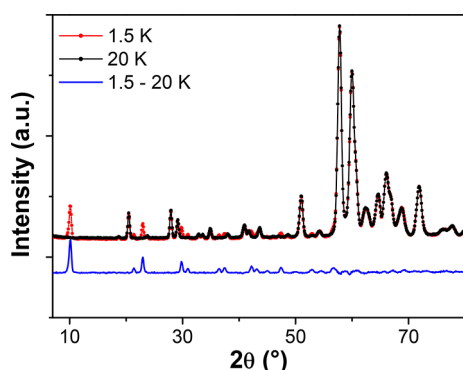
peak at 9 K indicating the establishment of an anti-ferromagnetic ordering. As shown in the inset of Figure 4, the inverse of the magnetic susceptibility measured at 3 T exhibits a small curvature in the paramagnetic domain. Thus, the data above  $T_N$  were fitted with a modified Curie–Weiss law:  $\chi = \chi_0 + C/(T - \theta_p)$ , where  $\chi_0$  is a temperature-independent term,  $C$  is the Curie constant, and  $\theta_p$  is the paramagnetic Curie temperature. From this fit, we obtain the following values:  $\chi_0 = 4.8 \times 10^{-4} \text{ emu} \cdot \text{mol}^{-1}$ ,  $C = 1.60 \text{ emu} \cdot \text{K} \cdot \text{mol}^{-1}$ , and  $\theta_p = -4 \text{ K}$ . The effective moment corresponding to the Curie constant is equal to  $3.58 \mu_B$ , which is very close to the value expected for the trivalent state of the  $\text{Nd}^{3+}$  ion ( $\mu(\text{Nd}^{3+}) = 3.62 \mu_B$ ). This also indicates that Ni is likely not or very weakly magnetic in this compound. The slightly negative paramagnetic Curie temperature (i.e.,  $-4 \text{ K}$ ), close to the Néel temperature, is coherent with local anti-ferromagnetic interactions between Nd atoms.



**Figure 5.** Field dependence of the magnetization curves at various temperatures between 1.8 and 15 K.

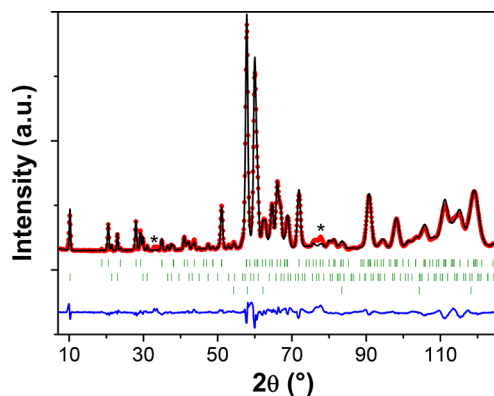
Figure 5 displays the isothermal magnetization curves at different temperatures between 1.8 and 15 K. Above  $T_N$ , the  $M(H)$  curve is linear as expected for a paramagnetic state. Below the Néel temperature, the curves exhibit a linear behavior at low fields, which is typical of an anti-ferromagnet. Besides, the curves show field-induced metamagnetic transitions at critical fields increasing with decreasing temperatures. For example, the critical field  $H_c$  at 1.8 and 7 K is  $\sim 4.8$  and  $2.75 \text{ T}$ , respectively (see Figure 5). This metamagnetic transition most likely results in a transition from an anti-ferromagnetic to a ferromagnetic state. The maximum value of the magnetization reaches  $\sim 2 \mu_B/\text{f.u.}$  at 7 T, but as the curve is clearly not completely saturated, one can expect a slightly higher saturated moment for Nd atoms.

**Magnetic Structure.** To know the type of anti-ferromagnetic order present in  $\text{NdNiMg}_{15}$  we performed neutron diffraction measurements above and below  $T_N$ , namely, between 20 and 1.8 K (see Figure 6). Above the Néel temperature, the Bragg peaks, indexed with the  $P4/nmm$  space group, only correspond to the nuclear contribution. At 20 K, the cell parameters are equal to  $a = 10.009(1) \text{ Å}$  and  $c = 7.731(1) \text{ Å}$ , and the refined crystal structure is in agreement with the results obtained from single-crystal measurements at room temperature. Some small extra peaks, not indexed with the  $P4/nmm$  space group, are accounting for the presence of



**Figure 6.** Neutron diffraction patterns of NdNiMg<sub>15</sub> recorded at 1.5 and 20 K with the difference pattern between both temperatures showing the magnetic contribution.

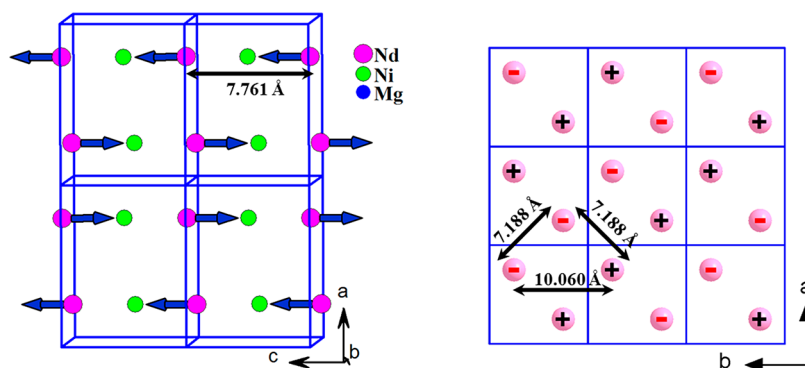
the residual secondary phases Mg, Mg<sub>2</sub>Ni, and NdMg<sub>12</sub>. Below 10 K, magnetic contributions start to grow at  $2\theta$  positions different from those of the nuclear contribution, as it can be seen in Figure 6. These magnetic peaks are indexed with the propagation vector  $\mathbf{k} = (1/2 \ 1/2 \ 0)$ , which is coherent with the anti-ferromagnetic state observed at  $T_N = 9$  K by magnetic measurements. This implies that the magnetic cell corresponds to the nuclear cell multiplied by 2 along both the  $a$ -axis and the  $b$ -axis ( $2a \times 2b \times c$ ). At 1.8 K, the best refinement of the difference diagram 1.8–20 K is unambiguously found for an anti-ferromagnetic model, where the magnetic moments are pointed along the  $c$ -axis. The refinement at 1.5 K and the corresponding magnetic structure are shown in Figures 7 and



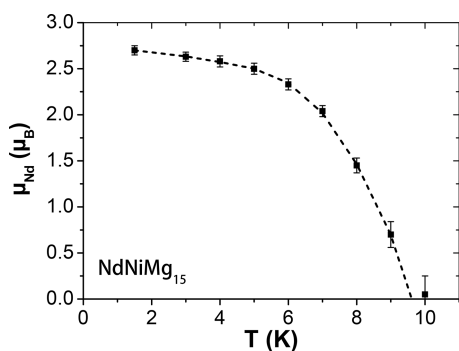
**Figure 7.** Rietveld refinement of the magnetic structure of NdNiMg<sub>15</sub> from the neutron diffraction pattern at 1.5 K. The two first lines of Bragg position ticks correspond to the nuclear and magnetic phases of NdNiMg<sub>15</sub>, respectively, and the third one corresponds to residual Mg phase. The symbol \* corresponds to the secondary phase Mg<sub>2</sub>Ni.

8. This magnetic structure consists in ferromagnetic chains along the  $c$ -axis of Nd atoms carrying moments, only separated by Ni atoms. These chains are coupled to each other ferromagnetically within planes perpendicular to the  $[110]$  direction and anti-ferromagnetically between these planes leading to a checkerboard pattern in the  $(ab)$  plane (Figure 8b). The Nd atoms form a puckered square lattice in the  $(a,b)$  plane with Nd–Nd distances of 7.19 and 10.06 Å. This latter distance corresponds to the diagonal of the square (see Figure 8). The “square lattices” of Nd atoms are stacked along the  $c$ -axis with Nd–Nd distances of 7.76 Å. In the  $(a,b)$  plane, anti-ferromagnetic couplings are observed between Nd atoms along

the  $[100]$  or  $[0-10]$  directions (Nd–Nd distance of 10.06 Å) and along the  $[110]$  directions for the Nd–Nd distance of 7.19 Å. Ferromagnetic couplings between Nd atoms are observed along  $[-110]$  directions (distance of 7.19 Å) and along the  $[001]$  direction (Nd–Nd distance of 7.76 Å). From the collinear magnetic structure, one can interpret the metamagnetic transition observed on the  $M(H)$  curve as a spin-flip transition. At the critical field value, half of the magnetic moments turn over along the direction of the field, which results in a ferromagnetic state. The magnetic moment value of Nd reaches  $\mu_{Nd} = 2.70(5) \mu_B$  at 1.5 K in consistency with the magnetization measurements  $M(H)$  at low temperatures. From the various neutron patterns recorded below  $T_N$ , we also refined  $\mu_{Nd}$  in function of the temperature as depicted in Figure 9. This thermal evolution confirms that the Néel temperature is  $\sim 10(1)$  K. No significant magnetic moment is found on the Ni site as expected from previous susceptibility measurement. However, the Néel temperature is remarkably high considering only the Nd sublattice with its high Nd–Nd distances within the crystal structure. It is known that the magnetic interactions between rare-earth atoms of Ruderman–Kittel–Kasuya–Yosida (RKKY) type are indirect and mediated by the spin polarization of conduction electrons induced by the localized 4f moments. This polarization, which has an oscillatory character, allows the interaction between localized 4f moments, and its absolute value decreases with increasing interatomic distance. The magnetic ordering temperature is also dependent on the density of states at the Fermi level. Among the other effects that can modulate this temperature one can cite the crystalline electric field<sup>41</sup> and the hybridization between the 5d orbitals of rare-earth and 3d orbitals of transition metal.<sup>42</sup> In NdNiMg<sub>15</sub>, with a Néel temperature of  $T_N = 9$  K, we observe very long Nd–Nd distances, equal to 7.76 and 7.19 Å between Nd atoms inside and between the chains, respectively. In the literature, only a few compounds having very long RE–RE distances (higher than 7 Å to achieve a 3D network of RE atoms) exhibit magnetic ordering temperature higher than 10 K. In the series RERhIn<sub>5</sub> with tetragonal symmetry, NdRhIn<sub>5</sub><sup>43,44</sup> is anti-ferromagnetic below  $T_N = 11$  K. The anisotropic structure is made of NdIn<sub>3</sub> and RhIn<sub>2</sub> layers with Nd–Nd distances equal to 4.63 and 6.55 Å in the NdIn<sub>3</sub> layers and to 7.50 Å between these layers. In the RE<sub>2</sub>Pt<sub>6</sub>Ga<sub>16</sub> series, a magnetic transition  $T_M \approx 13$  K is observed for RE = Nd.<sup>45</sup> These compounds crystallize with a hexagonal symmetry and are made of Pt<sub>2</sub>Ga<sub>14</sub> double layers and RE–Ga disordered layers. The RE–RE distances are nearly equal to 4.3 Å within the RE–Ga layers and to 8.7 Å between them. The last example is NdNiMg<sub>5</sub>,<sup>18</sup> which exhibits anti-ferromagnetic ordering below  $T_N = 12$  K, and its structure is made of Nd–Ni graphite-type and Mg layers. The Nd–Nd distances are equal to 4.48 and 5.47 Å inside the Nd–Ni layers and to 7.23 and 7.77 Å between these layers. It can be noticed that in all the examples, a strong structural anisotropy is observed. To show that the Nd–Nd distance is not the only criterion to take into account to estimate the magnetic transition temperature, the following examples are given. Nd<sub>2</sub>Zn<sub>17</sub><sup>46</sup> has Néel temperature close to 1.1 K, whereas the Nd–Nd distances are ranging from 4.35 to 6.84 Å. Moreover, in NdNi<sub>4</sub>Mg, Nd atoms form a tridimensional face-centered cubic (FCC) network with Nd–Nd distances of 5.02 Å, and no magnetic order is observed.<sup>47</sup> Band structure calculations are planned to better characterize the magnetic interactions in NdNiMg<sub>15</sub> and to see if the Ni atoms play a role in the interactions.



**Figure 8.** Magnetic structure of  $\text{NdNiMg}_{15}$ . (left) Four unit cells are represented in blue, and only Nd and Ni atoms are shown for clarity. (right) Representation along the  $c$ -axis. The (+) and (−) signs represent the magnetic moments pointing forward and backward along the  $c$ -axis, respectively.



**Figure 9.** Temperature dependence of the magnetic moment of Nd. The line is a guide.

## CONCLUSION

The search for new phases in the Mg-rich corner of the Nd–Ni–Mg was fruitful with the discovery of four phases, namely,  $\text{NdNiMg}_5$ ,<sup>18</sup>  $\text{Nd}_4\text{Ni}_8\text{Mg}_{80}$ ,<sup>20</sup>  $\text{Nd}_{16}\text{Ni}_{12}\text{Mg}_{96}$ ,<sup>21</sup> and the title compound  $\text{NdNiMg}_{15}$ . Its structure is an ordered variant of the  $\text{V}_{15}\text{Sb}_{18}$  structure, where a 3D network of Mg atoms and 1D chains of alternating Ni and Nd atoms are coexisting. Such arrangement with a 3D metallic network of Mg atoms and a covalent network of Ni and Nd atoms with lower dimensionality was previously observed in the other Mg-rich phases in the Nd–Ni–Mg system. Despite the very long Nd–Nd distances, higher than 7 Å, an anti-ferromagnetic ordering has been observed with  $T_N = 9$  K. The magnetic structure determined by neutron diffraction shows ferromagnetic coupling within the chains of Nd atoms, which are running along the  $c$ -axis. The chains are ferromagnetically coupled within planes perpendicular to the  $[110]$  direction, and these planes are anti-ferromagnetically coupled to neighboring planes forming a checkerboard-like magnetic structure. The metamagnetic transition observed on the  $M(H)$  curve can be interpreted as a spin-flip transition leading to an overall ferromagnetic structure.

## ASSOCIATED CONTENT

### Accession Codes

CCDC 1861893 contains the supplementary crystallographic data for this paper. These data can be obtained free of charge via [www.ccdc.cam.ac.uk/data\\_request/cif](http://www.ccdc.cam.ac.uk/data_request/cif), or by emailing [data\\_request@ccdc.cam.ac.uk](mailto:data_request@ccdc.cam.ac.uk), or by contacting The Cam-

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### Notes

The authors declare no competing financial interest.

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